

1 Title:

2 Clumped isotopologue constraints on the origin of methane at seafloor hot springs

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19 **Abstract**

20 Hot-spring fluids emanating from deep-sea vents hosted in unsedimented ultramafic and mafic rock commonly
21 contain high concentrations of methane. Multiple hypotheses have been proposed for the origin(s) of this me-
22 thane, ranging from synthesis via reduction of aqueous inorganic carbon (ΣCO_2) during active fluid circulation to
23 leaching of methane-rich fluid inclusions from plutonic rocks of the oceanic crust. To further resolve the pro-
24 cess(es) responsible for methane generation in these systems, we determined the relative abundances of several
25 methane isotopologues (including $^{13}\text{CH}_3\text{D}$, a “clumped” isotopologue containing two rare isotope substitutions) in
26 hot-spring source fluids sampled from four geochemically-distinct hydrothermal vent fields (Rainbow, Von
27 Damm, Lost City, and Lucky Strike).

28 Apparent equilibrium temperatures retrieved from methane clumped isotopologue analyses average 310_{-42}^{+53} °C,
29 with no apparent relation to the wide range of fluid temperatures (96 to 370 °C) and chemical compositions (pH,
30 $[\text{H}_2]$, $[\Sigma\text{CO}_2]$, $[\text{CH}_4]$) represented. Combined with very similar bulk stable isotope ratios ($^{13}\text{C}/^{12}\text{C}$ and D/H) of
31 methane across the suite of hydrothermal fluids, all available geochemical and isotopic data suggest a common
32 mechanism of methane generation at depth that is disconnected from active fluid circulation. Attainment of equi-
33 librium amongst methane isotopologues at temperatures of ca. 270 to 360 °C is compatible with the thermody-
34 namically-favorable reduction of CO_2 to CH_4 at temperatures at or below ca. 400 °C under redox conditions char-
35 acterizing intrusive rocks derived from sub-ridge melts. Collectively, the observations support a model where
36 methane-rich aqueous fluids, known to be trapped in rocks of the oceanic lithosphere, are liberated from host
37 rocks during hydrothermal circulation and perhaps represent the major source of methane venting with thermal
38 waters at unsedimented hydrothermal fields. The results also provide further evidence that water-rock reactions
39 occurring at temperatures lower than 200 °C do not contribute significantly to the quantities of methane venting at
40 mid-ocean ridge hot springs.

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1. INTRODUCTION

Dissolved methane (CH_4) is ubiquitous in hot-spring fluids emanating from submarine hydrothermal vents, and a potential carbon source for microbial communities living at and below the seafloor and in the water column. Constraining sources of carbon (C) and hydrogen (H) for the production of CH_4 , as well as the depths, temperatures, and timescales at which CH_4 is generated in these hydrothermal systems, is critical for understanding the origin of this form of hydrothermal carbon. The abundance and isotopic composition of CH_4 has been reported for fluids venting at unsedimented mid-ocean ridges and off-axis locations (e.g., Welhan, 1988b, Charlou et al., 2002, McCollom and Seewald, 2007, Proskurowski et al., 2008, Cannat et al., 2010, Charlou et al., 2010, Proskurowski, 2010, McDermott, 2015, and McDermott et al., 2015). In general, fluids that have interacted with ultramafic rock (peridotite and serpentinite) are enriched in CH_4 by one or more orders of magnitude relative to fluids that have reacted with mafic rock (basalt and gabbro) (Keir, 2010), although there are exceptions where high- CH_4 fluids emerge from mafic rock (Charlou et al., 2000).

Several distinct geochemical processes have been proposed to account for the presence of abiotic CH_4 in submarine hydrothermal fluids. Some have proposed that CH_4 is formed by reduction of aqueous inorganic carbon (i.e., $\sum\text{CO}_2$) during convective circulation of seawater-derived hydrothermal fluids in response to the highly reducing conditions that result from alteration of ultramafic rock (serpentinization) (Charlou et al., 2002; Proskurowski et al., 2008). Water-rock reactions during alteration of mafic rocks can also create sufficiently reducing conditions to reduce CO_2 to CH_4 without serpentinization (Shock, 1990; Seyfried and Ding, 1995), and is a possible source of the low but significant CH_4 observed in mafic-hosted systems (Keir, 2010). Others propose models involving entrapment and respeciation of mantle-derived CO_2 to CH_4 (and possibly to graphite) within plutonic (gabbroic) rocks of the oceanic crust (Kelley, 1996; Kelley, 1997; Kelley and Früh-Green, 1999), and subsequent extraction of the CH_4 -rich trapped fluids during hydrothermal circulation (McDermott et al., 2015). Leaching of basalt-hosted gas vesicles that contain CH_4 may also be a source of CH_4 in fluids venting at fast-spreading ridges such as the East Pacific Rise (Welhan and Craig, 1983; Welhan, 1988a).

To constrain the origin of CH_4 in unsedimented submarine hydrothermal systems, we determined the relative abundance of four of its stable isotopologues ($^{12}\text{CH}_4$, $^{13}\text{CH}_4$, $^{12}\text{CH}_3\text{D}$, and $^{13}\text{CH}_3\text{D}$, a doubly-substituted or “clumped” isotopologue) in nine fluid samples collected from four hydrothermal vent fields: Rainbow ($36^\circ13'48''$ N, $33^\circ54'09''$ W, Mid-Atlantic Ridge), Von Damm ($18^\circ22'36''$ N, $81^\circ47'54''$ W, Mid-Cayman Rise), Lost City ($30^\circ07'24''$ N, $42^\circ07'12''$ W, Mid-Atlantic Ridge), and Lucky Strike ($37^\circ17'30''$ N, $32^\circ16'42''$ W, Mid-Atlantic Ridge). These fields span a wide range of vent temperatures (96 to 370 °C), represent distinct geological settings, and are characterized by a wide range of fluid compositions.

Data presented in this study provide constraints on the sources of C and H, as well as temperature(s) associated with the formation or equilibration of C–H bonds in CH_4 carried by fluids of the hot-spring source. Bulk carbon- and hydrogen-isotope ratios ($^{13}\text{C}/^{12}\text{C}$ and D/H) encode signals related to the sources of C and H, respectively, as well as isotopic fractionations incurred during the synthesis of CH_4 . Complementary to such information, measurement of the CH_4 clumped isotopologue $^{13}\text{CH}_3\text{D}$ provides an independent estimate of the temperature at which the C–H bonds in CH_4 were formed or last equilibrated (Stolper et al., 2014; Wang et al., 2015). Constraining the temperatures at which CH_4 synthesis occurs within oceanic crust has direct implications for the distribution and availability of reduced carbon substrates and energy sources that may support a deep biosphere, as well as for the transfer of mantle-derived carbon to Earth’s surface.

85 Determination of temperatures from bulk carbon or hydrogen isotope ratios of CH₄ alone requires knowledge of
86 or assumptions regarding the isotopic composition of other species with which CH₄ has exchanged atoms (e.g.,
87 CO₂ or H₂O). In contrast, temperatures determined from the abundance of ¹³CH₃D do not require information
88 regarding such coexisting species. Thus, clumped isotopologue data in conjunction with bulk ¹³C/¹²C and D/H
89 isotope ratios of CH₄ can be used to constrain the isotopic compositions of C- and H-bearing species associated
90 with the CH₄ source when independent constraints are unavailable.¹ In the following discussion, we show how
91 clumped isotopologue temperatures of CH₄, together with bulk ¹³C/¹²C and D/H isotope ratios, fluid chemistry,
92 and thermodynamic considerations, indicate that CH₄ in un sedimented hydrothermal systems originates at tem-
93 peratures in excess of 270 °C and constrain possible environments of methane generation.

94 2. METHODS

95 2.1. Vent fluid samples

96 The fluid samples studied herein were collected by ROV *Jason II* using isobaric gas-tight samplers (Seewald et
97 al., 2002) during cruises to the Mid-Atlantic Ridge in 2008 (Reeves et al., 2014) and Mid-Cayman Rise in 2012
98 (McDermott et al., 2015). Subsamples of vent fluids extracted from the samplers were stored in pre-evacuated
99 serum vials sealed with blue butyl rubber stoppers that were preconditioned by boiling in 2 M NaOH for 2–4
100 hours and rinsed in deionized water. When necessary, sample aliquots in multiple serum vials were combined
101 (“pooled”) prior to purification to obtain enough CH₄ for clumped isotopologue analysis (>1 cm³ SATP). When
102 possible, aliquots from the same fluid sampler were used. In some cases, however, it was necessary to combine
103 aliquots from duplicate samples collected in separate samplers deployed in the same hydrothermal fluid during a
104 submersible dive (Table 1). Due to the exceedingly low concentration of dissolved CH₄ in ambient bottom sea-
105 water (<10⁻⁸ M, McDermott et al., 2015; Reeves et al., 2014) relative to concentrations in endmember vent fluids
106 (samples regressed to zero Mg content) (Table 2), inadvertent entrainment of seawater during fluid collection has
107 no measurable effect on the isotopic composition of CH₄ derived from vent fluids that was measured in this study.

108 2.2. Analytical techniques

109 Samples of CH₄ were purified via cryofocusing–preparative gas chromatography (Wang et al., 2015). The rela-
110 tive abundances of the methane stable isotopologues ¹²CH₄, ¹³CH₄, ¹²CH₃D, and ¹³CH₃D were measured using a
111 tunable infrared laser direct absorption spectroscopy technique described previously (Ono et al., 2014; Wang et
112 al., 2015). Due to the small amounts of CH₄ (ca. 1 cm³ STP) in samples analyzed as part of this study, a cold trap
113 system was employed to recover and recycle gas samples for re-analysis (Wang et al., 2015). A set of samples for
114 which isotopologue ratios had been previously determined was also re-measured using this recycling technique, to
115 verify accuracy (Supplementary Table 1).

116 The abundance of ¹³CH₃D relative to a random distribution of isotopes among the isotopologues (stochastic distri-
117 bution) is tracked using the metric Δ¹³CH₃D, which is defined as: Δ¹³CH₃D = ln *Q* (or nearly equivalently, *Q* – 1),
118 where *Q* is the reaction quotient of the isotope exchange reaction:



¹ This is analogous to using carbonate clumped isotope abundances to solve for the ¹⁸O/¹⁶O ratio of H₂O from which the car-
bonate precipitated. Readers are referred to Eiler (2007) for further information.

120 Values of $\Delta^{13}\text{CH}_3\text{D} > 0\text{‰}$ are used to calculate apparent equilibrium temperatures ($T_{13\text{D}}$) using the calibration of
121 Wang et al. (2015), which is based on quantum chemical predictions for CH_4 isotopologues in the gas phase and
122 anchored by measurements of methane samples heated in the presence of platinum catalyst at temperatures be-
123 tween 150 and 400 °C (Wang et al., 2015).

124 Bulk isotope values are reported herein using standard delta-notation, i.e., $\delta^{13}\text{C} = (^{13}\text{C}/^{12}\text{C})_{\text{sample}} / (^{13}\text{C}/^{12}\text{C})_{\text{VPDB}} - 1$,
125 and $\delta\text{D} = (\text{D}/\text{H})_{\text{sample}} / (\text{D}/\text{H})_{\text{VSMOW}} - 1$. The permil (‰) symbol represents multiplication by 10^{-3} ; hence, we have
126 omitted the factor of 1000 commonly seen in definitions of δ and other isotope values. The $\delta^{13}\text{C}$ and δD values
127 are calibrated against community reference gases NGS-1 and NGS-3 (Wang et al., 2015).

128 3. RESULTS

129 Results of stable carbon ($^{13}\text{C}/^{12}\text{C}$) and hydrogen (D/H) isotope ratio measurements are shown in Table 1. These
130 results are in general agreement with previously-published CH_4 isotopic data for these samples or systems
131 (Proskurowski et al., 2008; Charlou et al., 2010; Pester et al., 2012; McDermott et al., 2015). For fluids for which
132 direct comparisons to literature data are possible, the $\delta^{13}\text{C}$ values of CH_4 we report in Table 1 are typically $< 0.5\text{‰}$
133 different from those previously published. The largest deviation is a consistent 0.9‰ offset for all three samples
134 from Von Damm compared to the published report of McDermott et al. (2015) that likely stems from a difference
135 in calibration between the laboratories. Data for δD are sparse, but the δD value shown in Table 1 for Beehive
136 vent at Lost City is identical to that previously reported ($-127 \pm 6\text{‰}$; Proskurowski et al., 2008).

137 Similar isotopic values were observed across the different hydrothermal fields, ranging from -18‰ to -11‰ in
138 $\delta^{13}\text{C}$ and -127‰ to -98‰ in δD . Variation between vents in the same field (generally $< 1\text{‰}$ in both $\delta^{13}\text{C}$ and δD)
139 is significantly smaller than variation across different fields. The consistency of stable isotope data of CH_4 within
140 each field is added evidence for the interpretations previously drawn of conservative mixing of CH_4 between bot-
141 tom seawater and a single CH_4 -bearing endmember fluid at Rainbow (Charlou et al., 2002) and Von Damm
142 (McDermott et al., 2015). A common source fluid has also been suggested for Lucky Strike (Pester et al., 2012)
143 and Lost City (Seyfried et al., 2015) based on the compositions of fluids there.

144 Also shown in Table 1 are results of CH_4 clumped isotopologue analyses. All samples yielded values of $\Delta^{13}\text{CH}_3\text{D}$
145 $> 0\text{‰}$, from which apparent equilibrium temperatures can be derived (Fig. 1C and Table 1). The unweighted
146 mean of the $\Delta^{13}\text{CH}_3\text{D}$ values across all nine vent fluids studied is $1.57 \pm 0.28\text{‰}$ (standard deviation, $1s$), corre-
147 sponding to a $\Delta^{13}\text{CH}_3\text{D}$ temperature of 310^{+53}_{-42} °C (derived from projection of measured $\Delta^{13}\text{CH}_3\text{D}$ values onto the
148 green curve in Fig. 1C). Data for individual vent fluids are analytically indistinguishable from this narrow range
149 (Fig. 2B).

150 4. DISCUSSION

151 4.1. Closure temperatures for hydrogen exchange amongst CH_4 , H_2 , and H_2O in hydrothermal systems

152 The narrow range of measured $\Delta^{13}\text{CH}_3\text{D}$ values (averaging $1.57 \pm 0.28\text{‰}$, $1s$) and corresponding apparent equilib-
153 rium temperatures for Reaction 1 of 310^{+53}_{-42} °C contrasts with the wide range of fluid temperatures (96 to 370°C)
154 measured at the vents (Fig. 1, Table 2). Had the methane in these samples attained isotopologue equilibrium at
155 measured vent temperatures, $\Delta^{13}\text{CH}_3\text{D}$ values from 4.0 to 1.3‰ would be expected, respectively. The observed
156 range of clumped isotopologue data is much smaller than this predicted range (Fig. 1C), with $\Delta^{13}\text{CH}_3\text{D}$ tempera-

157 tures generally equal to or higher than fluid temperatures (Fig. 2B). The clumped isotopologue data indicate that
158 the bulk of CH₄ at the sites studied were either formed at or around 310 °C, or that CH₄ was generated elsewhere
159 in the hydrothermal system (perhaps at lower or higher temperatures) with subsequent establishment of equilibri-
160 um amongst the CH₄ isotopologues at 270 to 360 °C.

161 In addition to clumped isotopologues, bulk hydrogen isotope abundances also yield temperature constraints. Iso-
162 topic exchange of D/H amongst H₂O, H₂, and CH₄ can be described by the following equilibria:



166 Each of these reactions is characterized by a temperature-dependent equilibrium constant, such that differences in
167 the measured hydrogen isotopic ratios between CH₄-H₂O, H₂-H₂O, or H₂-CH₄ pairs can, in principle, be used to
168 calculate apparent equilibrium temperatures for the respective reactions. Seafloor hydrothermal fluids have δD
169 values of H₂O very close to 0‰ (i.e., seawater; Shanks et al., 1995; see Supplementary Table 2), such that differ-
170 ences in calculated isotopic temperatures between samples are almost entirely due to variations in δD values of
171 CH₄ or H₂.

172 Equilibrium fractionation factors between dissolved H₂ and CH₄ and liquid H₂O are shown in Fig. 3 (blue and red
173 curves, respectively). The fractionation factor for Reaction 3 (blue curves) was derived for each temperature by
174 multiplying experimental fractionation factors for H₂(g)/H₂O(g) of either Suess (1949) (thin solid line), Cerrai et
175 al. (1954) (thick solid line), or Bardo and Wolfsberg (1976) (dashed line), by that for H₂O(g)/H₂O(l) from Horita
176 and Wesolowski's (1994) calibration.² The fractionation factor for Reaction 2 (red curves) was then obtained by
177 multiplying the fractionation factors for Reaction 3 by the experimentally-calibrated CH₄(g)/H₂(g) fractionation
178 factor from Horibe and Craig (1995). Isotope effects of solvation were ignored because they are small (Muccitelli
179 and Wen, 1978; Bacsik et al., 2002).

180 Figure 3 shows a compilation of reported δD values of CH₄ at hydrothermal vent fields free of sediment influence.
181 All δD data of CH₄ from endmember fluids cluster around -110‰ ± 12‰ (1σ), and are consistent with CH₄ hav-
182 ing approached hydrogen isotopic equilibrium with seawater-like H₂O at temperatures in excess of 270 °C. An
183 upper temperature limit cannot be specified because of uncertainty in where the true D/H equilibrium fractiona-
184 tion between CH₄ and H₂O lies, and the low sensitivity of D/H thermometry of CH₄-H₂O (see Fig. 3). Tempera-
185 tures derived from Δ¹³CH₃D data are also in the range of 270 to 360 °C (shown by arrows), indicating that in sea-
186 floor hydrothermal systems, hydrogen exchange amongst CH₄ isotopologues is accompanied by or proceeds
187 through apparent hydrogen exchange between CH₄ and H₂O. Stated another way, Reaction 1 probably does not
188 proceed as an elementary reaction in nature. Instead, Reaction 1 likely evolves towards a state of equilibrium in-
189 directly as a consequence of other isotope exchange reactions such as that between CH₄ and H₂O (Reaction 2).
190 Reactions 1 and 2 therefore appear to proceed at similar rates, and closure temperatures³ for both reactions are

² Rolston et al. (1976) performed an experimental calibration of H₂(g)/H₂O(l) from 0 to 100 °C. The Rolston et al. curve is not shown in Fig. 3 because it covers only a portion of the temperature range of interest, but would plot very close to the dashed blue curve representing Bardo and Wolfsberg (1976) × Horita and Wesolowski (1994).

³ The *closure temperature* is the temperature below which isotopic values become “frozen” over observable time. Closure temperature is a function of cooling rate.

191 between 270 and 360 °C in hydrothermal systems. Below the closure temperature, processes that break and form
192 C–H bonds in CH₄ are slower than rates characterizing subsurface cooling of hydrothermal fluids.

193 Isotope data for H₂ are also shown in Fig. 3. Unlike for CH₄, H₂ is characterized by large spread in δD values
194 across sites (–700‰ to –330‰). Values for δD of H₂ strongly vary with measured vent temperature, with
195 endmember fluids plotting very near the equilibrium curves. This indicates that H₂ exchanges H with H₂O at rates
196 that keep pace with cooling of sub-seafloor fluids during ascent, and that the closure temperature for Reaction 3 is
197 much lower than that for CH₄–H₂O. An estimation based on data from low-temperature vents places the closure
198 temperature for abiotic H₂–H₂O exchange between 70 and 110 °C (gray arrows), albeit with large uncertainty.
199 Rapid metabolic cycling of H₂ by microorganisms living at the vents may enable H₂–H₂O exchange to occur
200 down to even lower temperatures (Proskurowski et al., 2006; Kawagucci et al., 2010).

201 The data also yield constraints on the kinetics of H₂–CH₄ exchange (Reaction 4). Figure 3 again shows that δD
202 values of H₂ plot very close to H₂–H₂O equilibrium while δD values of CH₄ plot away from the CH₄–H₂O line
203 below ~300 °C. This indicates that the closure temperature for H₂–CH₄ exchange must be high, probably not
204 much lower than 270 °C. If H₂–CH₄ were to actually close at a low temperature (100 °C for example), δD values
205 of CH₄ would instead plot close to the red curves because CH₄ would always be exchanging H with a population
206 of H₂ molecules that has a fixed δD value at each temperature (because of rapid exchange kinetics for H₂–H₂O
207 and a near-infinite pool of H in H₂O) shown by the blue curves. Indeed, experiments conducted with gas-phase
208 reactants between 25 and 400 °C show that rates for Reaction 3 (H₂–H₂O) are 20 to 100 times faster than that of
209 Reaction 4 (H₂–CH₄) given identical partial pressures of H₂ (Lécluse and Robert, 1994). It is possible that the
210 apparent exchange of hydrogen between CH₄ and H₂O at ~300 °C discussed in preceding paragraphs actually pro-
211 ceeds through H₂ (i.e., both CH₄ and H₂O undergo D/H exchange with H₂, while not interacting directly with each
212 other). If this is true, higher concentrations of H₂ would not only result in faster rates of H₂–CH₄ exchange, but
213 would also accelerate exchange between CH₄ and H₂O assuming rates of forward and backwards reactions in Re-
214 action 4 are first-order in both [H₂] and [CH₄].

215 The temperature-dependence of D/H exchange rates between CH₄, H₂, and liquid H₂O has been the subject of lim-
216 ited experimental study. Figure 4 summarizes the available experimental constraints for H₂–H₂O and CH₄–H₂O
217 exchange. Experiments with H₂ in the presence of liquid H₂O showed no observable exchange over 40 days at
218 26 °C (Campbell et al., 2009), but yielded half-exchange timescales as short as 10 min at 225 °C (Hall et al.,
219 1934; Lyon and Hulston, 1984). (The 225 °C experiments were conducted in the presence of potentially-catalytic
220 metal surfaces, and may therefore overestimate rates of exchange compared to nature.) Lécluse and Robert
221 (1994) reported second-order rate coefficients for isotope exchange between H₂ and H₂O in gas phase from 25 to
222 435 °C. To estimate exchange kinetics in liquid phase, we multiplied these second-order rate coefficients by the
223 vapor pressure of pure H₂O at each temperature, calculated at 1000 bar using a subroutine of the R package
224 CHNOSZ (Dick et al., 2008) from data of Johnson et al. (1992) and Wagner and Pruß (2002). The plotted rela-
225 tionship (the blue curve in Fig. 4) is consistent with calculations by Lin et al. (2005) who used data from the same
226 study, but experimental investigation is necessary to confirm that the kinetics of exchange for H₂ dissolved in wa-
227 ter are accurately represented.

228 Experiments on exchange between CH₄ and liquid H₂O are even scarcer, but available data suggest that half-
229 exchange timescales are on the order of 10 to 100 years at 300 °C, and ~10,000 years at 240 °C (Fig. 4; Supple-
230 mentary Fig. 1; Koeppe, 1978; Reeves et al., 2012). These numbers have large uncertainties because data collected
231 at very low extents of exchange must be extrapolated to obtain rate estimates. It is also not known how factors

232 such as pH, redox state, minerals, or concentrations of sulfur, H₂, or carbon species might affect hydrogen ex-
233 change rates. Nevertheless, experiments and observations demonstrate that CH₄-H₂O exchange is at least several
234 orders of magnitude slower than H₂-H₂O exchange.

235 4.2. Constraints on the origin of methane at seafloor hot springs

236 4.2.1. Evaluating potential microbial and thermogenic sources at unsedimented sites

237 A remaining question is whether CH₄ from the four vent fields we studied might be (i) thermogenic methane gen-
238 erated at >270 °C, or (ii) thermogenic or microbial methane generated during recharge at lower temperatures and
239 later heated during hydrothermal circulation through deeper hotter root zones. Bulk δD or Δ¹³CH₃D data may not
240 answer the second question since all C-H (and C-D) bonds rearrange upon heating to >270 °C, and thus C-H
241 bonds of CH₄ carry no information prior to heating above (and cooling below) ~270 °C. However, other data help
242 to exclude microbial and thermogenic sources as significant.

243 Bulk δ¹³C values of CH₄ fall within a relatively narrow range (-18‰ to -11‰) across fluids from all four un-
244 sedimented hydrothermal fields studied (Fig. 1A). Because the four sites we studied lack appreciable sedimenta-
245 tion, sedimentary organic carbon from which thermogenic hydrocarbons can be generated is in scarce supply
246 (Welhan, 1988b; Reeves et al., 2014). Dissolved organic carbon in recharging seawater is a potential source of
247 thermogenic methane, but can only account for ~40 μM (Sharp et al., 1995). In contrast, production of thermo-
248 genic methane by decomposition of organic matter at elevated temperatures is well documented from sediment-
249 influenced vent sites such as Guaymas Basin (Welhan and Lupton, 1987; Simoneit et al., 1988) and Middle Val-
250 ley (Cruse and Seewald, 2006). Methane from those sites are characterized by their low δ¹³C values (-40 to
251 -55‰) and high C₂₊ alkane concentrations (typical C₁/C₂ ratios <300), which are very different from the unsemi-
252 mented fields (δ¹³C > -24‰ and C₁/C₂ >1000) (McCollom and Seewald, 2007).

253 Methane from a CH₄-rich (~60 mM; Reeves et al., 2014) endmember fluid venting at 299 °C in Guaymas Basin
254 carried a Δ¹³CH₃D temperature of 326⁺¹⁷⁰₋₉₅ °C (95% confidence interval) (Wang et al., 2015). This sample also
255 had a δD value of CH₄ (-106‰) that is again consistent with hydrogen isotope equilibrium at ca. 300 °C with wa-
256 ter of VSMOW-like composition (large yellow octagon in Fig. 5).⁴ The similarity in δD and Δ¹³CH₃D values of
257 CH₄ between the sediment-influenced fluids and those from unsedimented fields, despite very different δ¹³C of
258 CH₄ (-44‰) and C₂₊ concentrations (McCollom and Seewald, 2007), suggests that at Guaymas Basin and other
259 high-temperature sediment-influenced fields (Kawagucci et al., 2013; Douglas et al., 2017), hydrogen isotopes of
260 CH₄ have also been reset by exchange with H₂O (and/or with H₂). We note that thermogenic natural gases gener-
261 ated at an early stage of kerogen maturation probably inherit C-H bonds from precursor organic molecules that
262 have undergone hydrogen exchange with other organics (instead of with water) in water-poor source rocks prior
263 to natural gas generation (Stolper et al., 2014; Wang, 2017). The latter CH₄ gases evolve towards apparent D/H
264 equilibrium with water with increasing maturity (shown by the brown arrow in Fig. 5) (Clayton, 2003), perhaps
265 by inheritance of H from methyl moieties that have previously exchanged with water (Hoering, 1984; Smith et al.,
266 1985; Lewan, 1997; Schimmelmann et al., 1999; Schimmelmann et al., 2001; Seewald, 2003; Lis et al., 2006;

⁴ Note added in revision: Methane clumped isotope data have now been reported from two additional seafloor hot-spring flu-
ids. The samples are from the Main Endeavour field on the Juan de Fuca Ridge (Douglas et al., 2017). These fluids have
elevated concentrations of CH₄ (millimolar) and low δ¹³C values (ca. -50‰), indicating the presence of buried and heated
sedimentary organic matter that contributes thermogenic CH₄ (Lilley et al., 1993). The CH₄ is relatively enriched in deuterium
(ca. -100‰) and carry high clumped isotope temperatures (~300 to 360 °C). These values are very similar to the Guay-
mas Basin sample from Wang et al. (2015) (see Fig. 5), providing additional evidence that apparent hydrogen-exchange oc-
curs in high-temperature fluids from sediment-influenced fields.

267 Schimmelmann et al., 2006; Reeves et al., 2012; Wang, 2017), or in the case of the hot Guaymas Basin fluids,
268 possible direct exchange of hydrogen between CH₄ and H₂O.

269 Microbial methane is generally characterized by low δ¹³C values (<-60‰) due to carbon isotope fractionation
270 during microbial methanogenesis. The methane δ¹³C data alone, however, do not unambiguously exclude contri-
271 butions of microbial methanogenesis, because high methane δ¹³C values could be a result of near-quantitative
272 conversion of ΣCO₂ to CH₄, particularly under ΣCO₂-limited conditions such as those present at Lost City
273 (Brazelton et al., 2006; Bradley and Summons, 2010). Moreover, Takai et al. (2008) reported that carbon isotope
274 fractionation by methanogens utilizing ΣCO₂ at >120 °C becomes very small (<12‰) under high H₂ partial pres-
275 sure (150 bar), suggesting that microbially-derived CH₄ might have similar δ¹³C values as those observed in un-
276 sedimented vent fluids (ca. 10 to 20‰ lower than seawater ΣCO₂, Fig. 1A). However, radiocarbon (¹⁴C) abun-
277 dances in CH₄ from Lost City and Von Damm are very low [fraction modern (*F_m*) averaging 0.004–0.006, near
278 the limit of detection (*F_m* ~ 0.003)] (Proskurowski et al., 2008; McDermott et al., 2015), whereas ¹⁴C contents of
279 endmember ΣCO₂ at Von Damm are ~5× higher (McDermott et al., 2015). Had CH₄ been derived from reduction
280 of ΣCO₂, the younger ¹⁴C age of the ΣCO₂ would have been transferred to the CH₄ product. McDermott et al.
281 (2015) further showed that ΣCO₂ in the vent fluids at Von Damm is likely seawater-derived, because both con-
282 centrations and δ¹³C values of endmember ΣCO₂ match those of local bottom seawater. The conservation of
283 ΣCO₂ during convective circulation at Von Damm excludes any process—microbial or otherwise—that converts
284 ΣCO₂ in recharging seawater at Von Damm to CH₄ despite high energetic favorability for CH₄ synthesis at *in situ*
285 conditions (Fig. 6D). The similarity in isotopic data and C₁/C₂ concentration values across the multiple vent
286 fields studied here suggests that processes responsible for CH₄ generation at Von Damm are also occurring at the
287 other unsedimented sites, even if the predominant influences on bulk fluid chemistry differ between sites. There-
288 fore, the data described above support the idea that the endmember-derived CH₄ in the studied hydrothermal flu-
289 ids is of dominantly abiotic origin (e.g., Charlou et al., 2002, McDermott et al., 2015, Proskurowski et al., 2008,
290 and Welhan, 1988), and that the contribution of thermogenic or microbial processes to the CH₄ content of the
291 endmember source component of fluids venting at each site is limited or insignificant.

292 4.2.2. High temperature origin of CH₄ at Lost City

293 The isotopologue data discussed above indicate that at Rainbow, Lucky Strike, Von Damm, and Lost City, CH₄
294 experienced temperatures in excess of 310⁺⁵³₋₄₂ °C at least once in its lifetime. The same processes responsible for
295 the rupturing and healing of C–H bonds by which CH₄ isotopologues attain equilibrium may also drive concentra-
296 tions of CH₄ to thermodynamic equilibrium. This is because the rate-limiting step for conversion of ΣCO₂
297 to/from CH₄ in hydrothermal fluids is thought to be the conversion of methanol (CH₃OH) to/from CH₄ (i.e., the
298 formation/breakage of C–H bonds in CH₄) (Seewald et al., 2006; McCollom and Seewald, 2007; Reeves, 2010).
299 Thus, Δ¹³CH₃D and δD values of CH₄ likely record the most recent temperature at which CH₄ synthesis was both
300 thermodynamically favorable and kinetically facile. This means that either net synthesis of CH₄ occurred at or
301 above ca. 300 °C, or that CH₄ already existed and re-equilibrated at this temperature under conditions where CH₄
302 synthesis would have proceeded had there been ΣCO₂ available.

303 The Δ¹³CH₃D temperature of 270⁺¹⁰⁴₋₆₈ °C we obtained for the Beehive vent fluid at Lost City argues for a much
304 higher temperature of last exchange for the C–H bonds in methane than Proskurowski et al.'s (2006) suggestion
305 of 110 to 150 °C. Their conclusion was based on hydrogen-isotope geothermometry of H₂–CH₄, and was prem-
306 ised on the assumption that δD values of H₂ and CH₄ in the Lost City vent fluids reflect isotope equilibrium be-
307 tween these species. However, as discussed earlier, rapid H₂–H₂O equilibration at temperatures lower than the

308 closure (quench) temperatures for CH₄–H₂O or H₂–CH₄ exchange will shift δD values for H₂ progressively lower,
309 while δD values of CH₄ remain unchanged. The fact that apparent temperatures derived from H₂–CH₄ will always
310 be close to that from H₂–H₂O is a coincidence that results from equilibrium δD values for CH₄ being so poorly
311 sensitive to temperature in comparison to those of H₂ (Fig. 3). Our Δ¹³CH₃D data therefore suggest that at the
312 most recent point at which CH₄ at Lost City was at isotopic equilibrium, surrounding temperatures were in excess
313 of 200 °C, an assertion that is also supported by the bulk δD values of CH₄. The high methane clumped isotopo-
314 logue temperature does not imply that the relatively cool fluids that vent at the surface today at Lost City experi-
315 enced such high temperatures, because CH₄ may have been formed elsewhere and been entrained into cooler cir-
316 culating fluids prior to their ascent to the seafloor (see Sec. 4.2.3).

317 4.2.3. *A deep origin of methane disconnected from actively circulating fluids*

318 The reduction of ΣCO₂ to CH₄ under hydrothermal conditions can be described by:



320 Reduction of CO₂ is favored at temperatures below ~300 °C under redox conditions near the pyrite-pyrrhotite-
321 magnetite (PPM) redox buffer, which approximates typical conditions for hydrothermal vent systems (Fig. 2A)
322 (Shock, 1992; Seyfried and Ding, 1995). At sites with ultramafic rocks, reactions involved in the hydration of
323 mafic minerals (serpentinization) can produce highly reducing conditions (H₂ fugacity more than 100-fold higher
324 than PPM, Fig. 2). Serpentinization occurs at temperatures below 400 °C, where ferrous iron-bearing olivine and
325 orthopyroxene minerals become unstable, with the highest H₂ fugacities produced at temperatures of 300 to
326 330 °C during serpentinization of peridotite under low water-to-rock ratios (Sleep et al., 2004; Klein et al., 2009;
327 Klein et al., 2013). This range of temperatures is generally consistent with the closure temperatures indicated by
328 our Δ¹³CH₃D data, suggesting that formation of CH₄ probably did not occur at temperatures much higher than ca.
329 400 °C (Fig. 2B). However, while ΣCO₂ reduction (Reaction 5) is thermodynamically favorable under conditions
330 encountered at ultramafic-hosted vents (Fig. 6D), the rate of the reaction is slow even at high temperatures
331 (~300 °C). Numerous long-term experiments conducted between 177 and 325 °C with ¹³C-labeled ΣCO₂ show
332 that production of CH₄ is kinetically-hindered in the absence of metal catalysts such as native iron (McCollom
333 and Seewald, 2001; McCollom and Seewald, 2003; Seewald et al., 2006; Reeves, 2010; McCollom, 2016;
334 Grozeva et al., 2017). In most of these experiments, uncatalyzed rates of ΣCO₂ conversion of less than a few per-
335 cent per year are observed. Considering that estimated residence times of fluids in the high-temperature reaction
336 zone (>200 °C) of several mid-ocean ridge systems is years to decades (Kadko, 1996; Fisher, 2003), reduction of
337 ΣCO₂ to CH₄ within actively-circulating hydrothermal fluids may be quite limited. Indeed, the isotopic composi-
338 tion and abundance of aqueous carbon species at the Von Damm vent field indicate that ΣCO₂ reduction to CH₄ is
339 not occurring on the timescale of convective hydrothermal circulation there (see Sec. 4.2.1 and McDermott et al.,
340 2015). We suggest that CH₄ in the hot-spring fluids studied here originates not within actively-circulating fluids,
341 but instead formed elsewhere and was entrained into the fluids prior to venting, as was suggested for Von Damm
342 by McDermott et al. (2015).

343 Concentrations of CH₄ for the four hydrothermal fields investigated in this study range from 0.86 to 2.81 mM
344 (Fig. 6B). Millimolar concentrations are typical of many ultramafic-hosted mid-ocean ridge hydrothermal fields,
345 whereas basalt-hosted fields tend to have lower CH₄ contents (~0.05 to 1 mM; McCollom and Seewald, 2007;
346 Keir, 2010). While CH₄ concentrations vary less than four-fold (from 0.9 to 2.8 mM), the studied fluids show a
347 wide range of pH (3.3 to 10.2), ΣCO₂ (<0.2 to 110 mM) and H₂ (0.03 to 18.2 mM) concentrations (Table 2).

348 The concentration of dissolved H₂ is high and varies from 10.4 to 18.2 mM in endmember fluids from the Rain-
349 bow, Von Damm, and Lost City fields, whereas fluids from Lucky Strike have concentrations of H₂ that are ap-
350 proximately three orders of magnitude lower (34–63 μM, Fig. 6A). At Rainbow, Von Damm, and Lost City, ser-
351 pentinization of ultramafic rock in subsurface reaction zones (and concomitant H₂ production) is thought to be a
352 major control on fluid compositions (Kelley et al., 2001; Charlou et al., 2002; McDermott et al., 2015). In con-
353 trast, the Lucky Strike field is hosted in mafic rocks, and vent fluids there encounter much more oxidizing condi-
354 tions (Charlou et al., 2000; Pester et al., 2012). At Lucky Strike, synthesis of CH₄ within the low-H₂ endmember
355 fluids is thermodynamically unfavorable at *in situ* temperatures (270–292 °C; Table 2 and Fig. 1B), and becomes
356 even more unfavorable with increasing temperature (Fig. 1A). In all other vent fluids studied here, a thermody-
357 namic drive for CH₄ synthesis is present at varying magnitudes (Fig. 6D). Thus, apparent decoupling of CH₄ con-
358 centrations from thermodynamic drives dictated by vent fluid chemistry (ΣCO₂, H₂ and pH) and temperature sug-
359 gests that CH₄ is generated independently at depth, and that conditions under which CH₄ forms are different from
360 those observed at venting.

361 Endmember fluids from Rainbow, Von Damm, and Lucky Strike contain 2 to 50 times as much total carbon as is
362 in bottom seawater (~2.2 mM; McDermott et al., 2015; Reeves et al., 2014), such that ΣCO₂ in recharging sea-
363 water cannot be the sole source of carbon to venting fluids. The highly-alkaline Lost City fluid (pH 10.2) con-
364 tains very low amounts of ΣCO₂ (<0.18 mmol/kg), the majority of which is likely derived from seawater entrain-
365 ment during sample collection (Reeves et al., 2014). At Lost City, Proskurowski et al. (2008) suggested that CH₄
366 forms via ΣCO₂ reduction within the circulating fluids. This model is problematic because it requires the addition
367 of mantle-derived CO₂ to the fluid and immediate reduction to CH₄ thereafter before ΣCO₂ can be removed by
368 carbonate precipitation. While sluggish kinetics characterize reduction of ΣCO₂ to CH₄, carbonate precipitation
369 proceeds much more rapidly at temperatures experienced by the circulating fluids (Kelemen et al., 2011; Grozeva
370 et al., 2017). (We note that it is possible to generate CH₄ directly from carbonate minerals, for example via hy-
371 drogenation reactions, as demonstrated by Giardini et al. (1968) and Yoshida et al. (1999), among others. How-
372 ever, those experiments were carried out with metal (e.g., Co, Ni and Cu) carbonate, the reactions involved are
373 remarkably slow below 400 °C, and their relevance to natural systems is doubtful.) This difference in kinetics
374 means that barring some unidentified catalytic process, ΣCO₂ reduction within the circulating hydrothermal flu-
375 id—and particularly at the low temperatures postulated by Proskurowski et al. (2006)—does not explain the CH₄
376 venting from Lost City because ΣCO₂ reduction cannot occur in the absence of there being any ΣCO₂ to reduce.

377 The clumped isotopologue temperature obtained for the Lost City methane (270⁺¹⁰⁴₋₆₈ °C) overlaps with or is
378 slightly higher than temperature estimates from heat balance considerations, δ¹⁸O values, and alkane-alkene and
379 mineral-fluid equilibria that all suggest that fluids beneath Lost City experienced temperatures as high as 200 to
380 250 °C (Allen and Seyfried, 2004; Foustoukos et al., 2008; Reeves et al., 2012; Seyfried et al., 2015). Agreement
381 between temperatures derived from Δ¹³CH₃D and these other geothermometers could be a matter of circumstance
382 (e.g., if all have similar closure temperatures), however, and does not necessarily mean that CH₄ formed and
383 equilibrated its C–H bonds within the same fluids that the other equilibria are recording. Instead, CH₄ may have
384 formed and attained isotopologue equilibrium independently of (outside) the circulating fluids, and later was en-
385 trained into them prior to venting. For example, a component of the hydrothermal fluid may have percolated very
386 deeply via meandering flow paths (Titarenko and McCaig, 2016), taking significantly longer to reach seafloor
387 vents and seeing higher temperatures than the remainder of the vent fluid (Hasenclever et al., 2014). If this near-
388 ly-stagnant “long-path” fluid was also CH₄-laden, mixing of a minute proportion of this fluid with a CH₄-poor
389 “short-path” hydrothermal fluid could be the source of the up to millimolar quantities of CH₄ emanating from sed-
390 iment-free seafloor vent systems.

391 Fluid inclusions record the widespread occurrence and composition of such a CH₄-rich aqueous fluid within the
392 sub-oceanic ridge lithosphere. Kelley (1996; 1997) and Kelley and Früh-Green (1999) documented several types
393 of volatile-rich inclusions hosted in plutonic rocks (gabbros) recovered from the slow-spreading Southwest Indian
394 and Mid-Atlantic Ridges by several Ocean Drilling Program (ODP) expeditions. They noted a common type of
395 inclusion occurring along healed microcracks in plagioclase grains (i.e., secondary inclusions) that contained up
396 to 47 mole percent CH₄ (with balance of H₂O), as well as possibly graphite and H₂. Temperatures indicated by
397 CO₂-CH₄ carbon isotope geothermometry (300–600 °C) and homogenization temperatures of the Southwest Indi-
398 an Ridge fluid inclusions (350–370 °C, corresponding to entrapment at *in situ* temperatures of ca. 400 °C) (Kelley
399 and Früh-Green, 1999) were interpreted to indicate formation of CH₄ during re-speciation of trapped magmatic
400 volatiles from CO₂ to CH₄ (± graphite) as the melt-derived host rocks cooled to below 400 °C at redox conditions
401 near the FMQ buffer (fayalite-magnetite-quartz, Fig. 2A; Kelley, 1996).

402 The δ¹³C values of CH₄ in vent fluids from unsedimented fields (–24 to –6‰, Supplementary Fig. 2A; Keir,
403 2010; McCollom and Seewald, 2007) are generally consistent with those determined by analyses of CH₄ in the
404 inclusions. Values for inclusions can be somewhat lower (–34 to –20‰; Kelley and Früh-Green, 1999), probably
405 due to contamination by relatively ¹³C-depleted thermogenic CH₄ released from background carbon sources dur-
406 ing heating of samples to decrepitate the inclusions. In both vent fluids and inclusions, CH₄ is generally more
407 ¹³C-depleted than mantle-derived CO₂ (–5‰; references in Fig. 1A caption).⁵ This presents a mass-balance issue
408 because a ¹³C-enriched component is apparently missing. Consistent with this, CH₄/³He ratios in vent fluids (see
409 Keir, 2010) indicate less-than-quantitative conversion (~0.2% to 50%) of mantle carbon to CH₄ (assuming mantle
410 C/³He is 1×10⁹, Marty and Tolstikhin, 1998). Precipitation of graphite from a CH₄-rich fluid entrapped in plutonic
411 rocks may explain both the missing carbon (McDermott et al., 2015) and the observed δ¹³C values (Luque et
412 al., 2012). Graphite precipitation in trapped fluids, suggested by Kelley (1996) and others, is consistent with
413 thermodynamic calculations that show that graphite can co-exist with CH₄ at ca. 400 °C under H₂O-poor condi-
414 tions and redox situated close to FMQ (French, 1966; Eugster and Skippen, 1967; Ohmoto and Kerrick, 1977;
415 Holloway, 1984; Früh-Green et al., 2004).

416 Concentrations of CH₄ in the fluid inclusions (up to 47 mole percent) from the Southwest Indian Ridge and from
417 other slow-spreading areas can be several orders of magnitude greater than those observed in corresponding vent
418 fluids (on the order of 1 mmol/kg, or 0.002 mole percent) (Kelley, 1996; Kelley, 1997). Mass-balance analysis
419 indicates that extraction of CH₄-rich fluids trapped in plutonic rocks can explain the observed CH₄ concentrations
420 at sediment-free mid-ocean ridge hydrothermal fields. Mixing curves plotted in Fig. 7 show that addition of less
421 than 0.1% of a CH₄-rich fluid of similar composition to those indicated by the inclusions (10 mole percent; *Fluid*
422 *2* in the figure) to a CH₄-poor circulating hydrothermal fluid (*Fluid 1*) is sufficient to match even the highest CH₄
423 concentrations seen in vent fluids. Assuming carbon contents ranging from 30 to 300 ppm in the gabbro (Kelley
424 and Früh-Green, 1999), water-to-rock ratios between 0.8 and 8 can explain CH₄ concentrations of up to
425 3 mmol/kg in vent fluids assuming all carbon in gabbro existed as leachable CH₄. These water-to-rock ratios are
426 consistent with constraints from Li, Rb, and Sr, which indicate ratios <<10 in many mid-ocean ridge hydrothermal
427 systems (Von Damm et al., 1985; Berndt et al., 1989); values of 0.4 to 6 and 2 to 4, respectively, were calculated
428 at Von Damm and Lost City for instance (Foustoukos et al., 2008; McDermott, 2015).

429 The above discussion shows that production of CH₄ within actively-circulating hydrothermal fluids is unlikely to
430 account for concentrations and isotopic signatures observed at several studied vent fields. Instead, our data are

⁵ Supplementary Fig. 2B shows that δ¹³C values of CH₄ and CO₂ in seafloor vent fluids are not correlated. This provides further evidence that CH₄ does not derive from CO₂ in actively-circulating fluids.

431 compatible with the idea that CH₄-rich fluids trapped in plutonic rocks are liberated during convective hydrother-
432 mal circulation, and that the fluids that the CH₄-rich inclusions represent could be the major source of CH₄ in hot-
433 spring fluids at sediment-free oceanic spreading centers. We speculate that these inclusions may have formed
434 during reactions between magmatic volatiles and mafic minerals during incipient percolation of seawater into
435 melt-derived rocks. Available data do not allow us to determine whether direct leaching of the fluid inclusions
436 themselves is the source of CH₄ to endmember hot-spring fluids, or whether both the vent fluid CH₄ and the fluids
437 cached in the inclusions have a common source (i.e., the inclusions simply record a prior passage through plutonic
438 rocks, via small fractures and/or mineral interstices, of the same slowly-moving CH₄-laden parent fluid as those
439 that contribute the CH₄ venting today); both are possible and indeed likely.

440 While only slow-spreading environments were investigated in this study, the same origin of CH₄ might apply at
441 sites on fast-spreading ridges such as the East Pacific Rise, particularly given the similar $\delta^{13}\text{C}$ values for CH₄
442 (Welhan and Craig, 1983). Concentrations of CH₄ (and C₂₊) in vent fluids there tend to be lower relative to sites
443 on slow-spreading ridges (Welhan, 1988b; Keir, 2010). Differences in axial structure and tectonism may account
444 for this pattern. At magma-poor slow-spreading ridges, extension is accommodated primarily by detachment
445 faulting, as opposed to magmatic emplacement of new crust that characterizes fast-spreading ridges (Buck et al.,
446 2005; Dunn, 2007). Low-angle, large-offset, and long-lived (>1 Myr) normal faults near vent fields at slow-
447 spreading ridges allow for fluid penetration deep into plutonic rocks of layer 3, enabling access to fresh gabbroic
448 material and/or inclusions to be leached (Kelley, 1996; Schroeder et al., 2002; Schlindwein and Schmid, 2016).
449 In contrast, in fast spreading environments such as the East Pacific Rise, shallow melt lenses at 1 to 2 km below
450 seafloor may limit the depth of circulation (e.g., Hasenclever et al., 2014, and references in Alt, 1995). Lucky
451 Strike, though hosted in mafic rock, is characterized by an unusually deep reaction zone (>3 km) and axial magma
452 chamber, as well as deeply penetrating fault reflectors (e.g. Pester et al., 2012, Escartin et al., 2015 and references
453 therein); thus, circulation there could lead to increased leaching of CH₄ relative to mafic-hosted systems in faster-
454 spreading settings. Efforts to define CH₄ origin will benefit from rigorous interrogation of factors governing fluid
455 flow and chemical kinetics in hydrothermally-influenced settings.

456 5. CONCLUSIONS

457 Measured abundances of methane isotopologues in fluids venting from diverse unsedimented mid-ocean ridge
458 hydrothermal systems are highly uniform, and yield last equilibration temperatures of ca. 300 °C for the C–H
459 bond. Taken in combination with geochemical and geologic observations and reaction rates determined in exper-
460 iments, the $\Delta^{13}\text{CH}_3\text{D}$ data establish that abiogenic reduction of ΣCO_2 (via e.g., FTT synthesis) at low temperatures
461 (<200 °C) is unlikely to be a significant source of methane over timescales characterizing convective hydrother-
462 mal circulation at oceanic spreading centers. Apparent decoupling between methane isotopologue abundances
463 and vent fluid chemistry points to a deep origin of CH₄ that is disconnected from active hydrothermal circulation.

464 We believe that the available data are best explained by a model where fluids rich in CH₄ (up to tens of mole per-
465 cent) form within plutonic rocks during re-speciation of magmatic carbon trapped in fluid inclusions or interstitial
466 spaces between mineral grains at temperatures at above 300 °C and under low water-to-rock ratios, and are later
467 leached into circulating hydrothermal fluids. Vent fluids with millimolar quantities of CH₄ therefore represent
468 mixing of a minute amount of a CH₄-rich fluid with a large volume of an actively-circulating, CH₄-poor fluid.
469 Proportions of mixing may be determined by the relative access that circulating fluids have to magmatic volatile-
470 bearing rocks of the plutonic foundation. Such a model could explain apparent relationships of CH₄ concentration
471 in vent fluids to tectonic setting and host rock lithology. It also explains the apparent lack of correlation between

472 CH₄ concentrations and thermodynamic drive for CH₄ synthesis calculated for chemical compositions and tem-
473 peratures of the venting fluids.

474 The new data also place constraints on the closure temperature of hydrogen exchange between methane and wa-
475 ter. The observation of sluggish or indiscernible exchange of H among methane isotopologues below ca. 300 °C
476 on timescales of ~10² years is relevant not only to the application of clumped isotope measurements as a novel
477 geothermometer, but also provides information about the stability of the C–H bond in hydrocarbons in nature.
478 Given the increasing appreciation of hydrocarbon-water-mineral interactions in economically important settings
479 (Seewald, 2003), insights of this nature may find utility in studies of the origin and composition of aqueous and
480 organic fluids in the Earth’s subsurface.

481

482

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492

493 **Supplementary Material** is included with this submission.

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496

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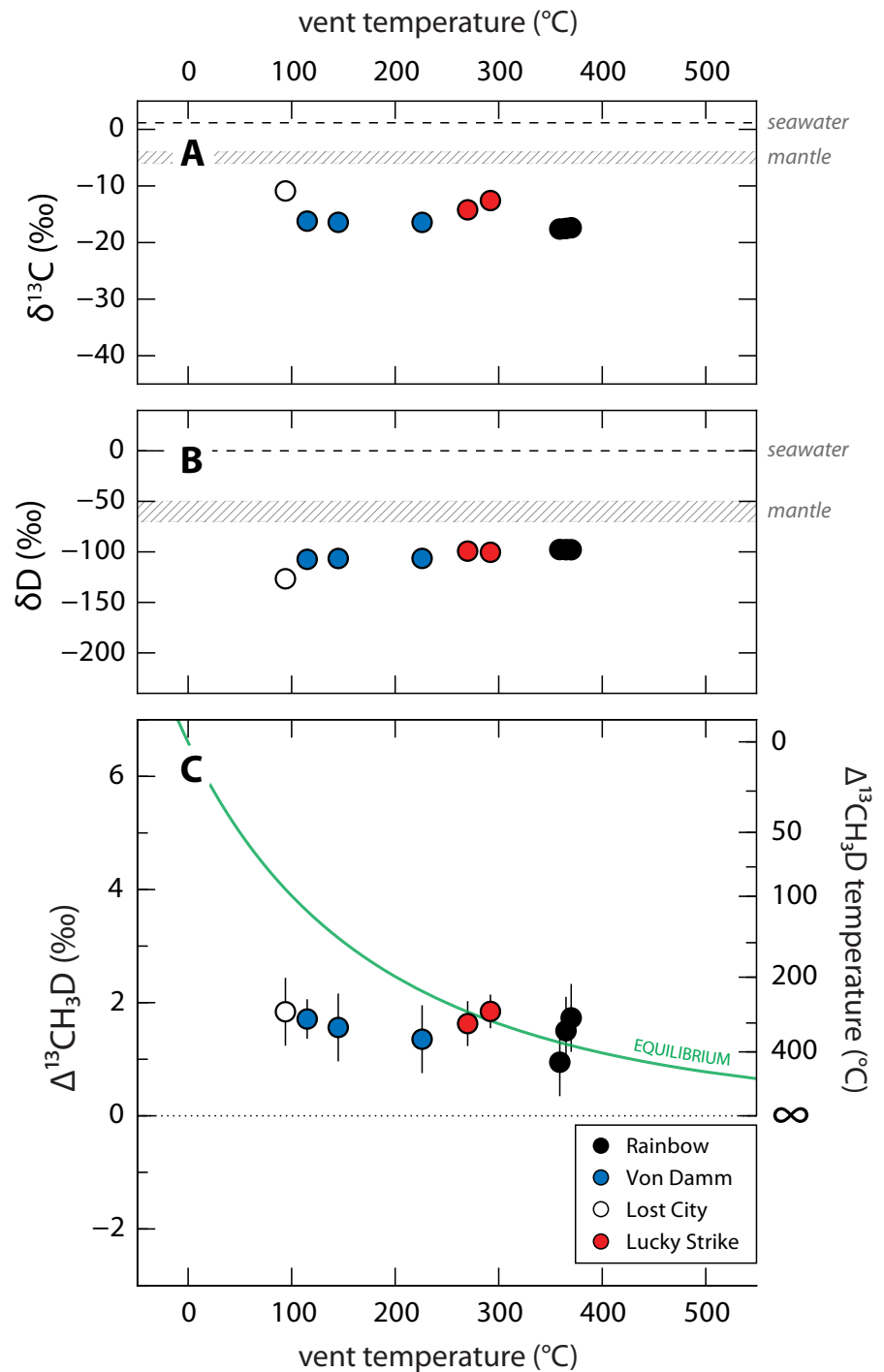
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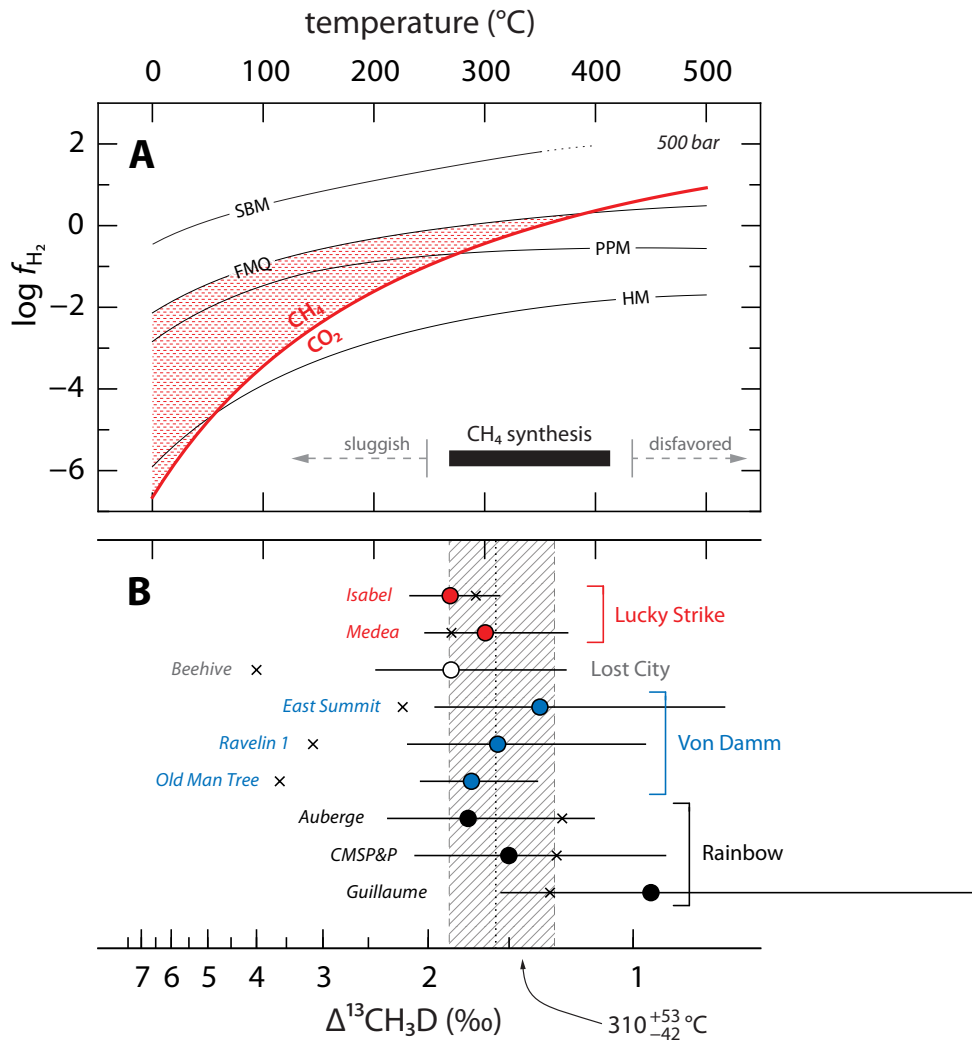
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782 **Fig. 1.** Comparison of (A) $\delta^{13}\text{C}$,
783 (B) δD , and (C) $\Delta^{13}\text{CH}_3\text{D}$ values of
784 methane across vent sites. Data
785 and error bars (95% confidence
786 interval) are from Table 1. In all
787 panels, data are plotted against
788 measured vent temperature (Table
789 2). The isotopic compositions of
790 inorganic carbon (A) and hydrogen
791 (B) in seawater and in the mantle
792 are shown (Javoy et al., 1986;
793 Blank et al., 1993; Clog et al.,
794 2013). In (C), the green line repre-
795 sents the clumped isotopologue
796 composition at equilibrium. The
797 $\Delta^{13}\text{CH}_3\text{D}$ temperature scale corre-
798 sponds to the calibration from
799 Wang et al. (2015).

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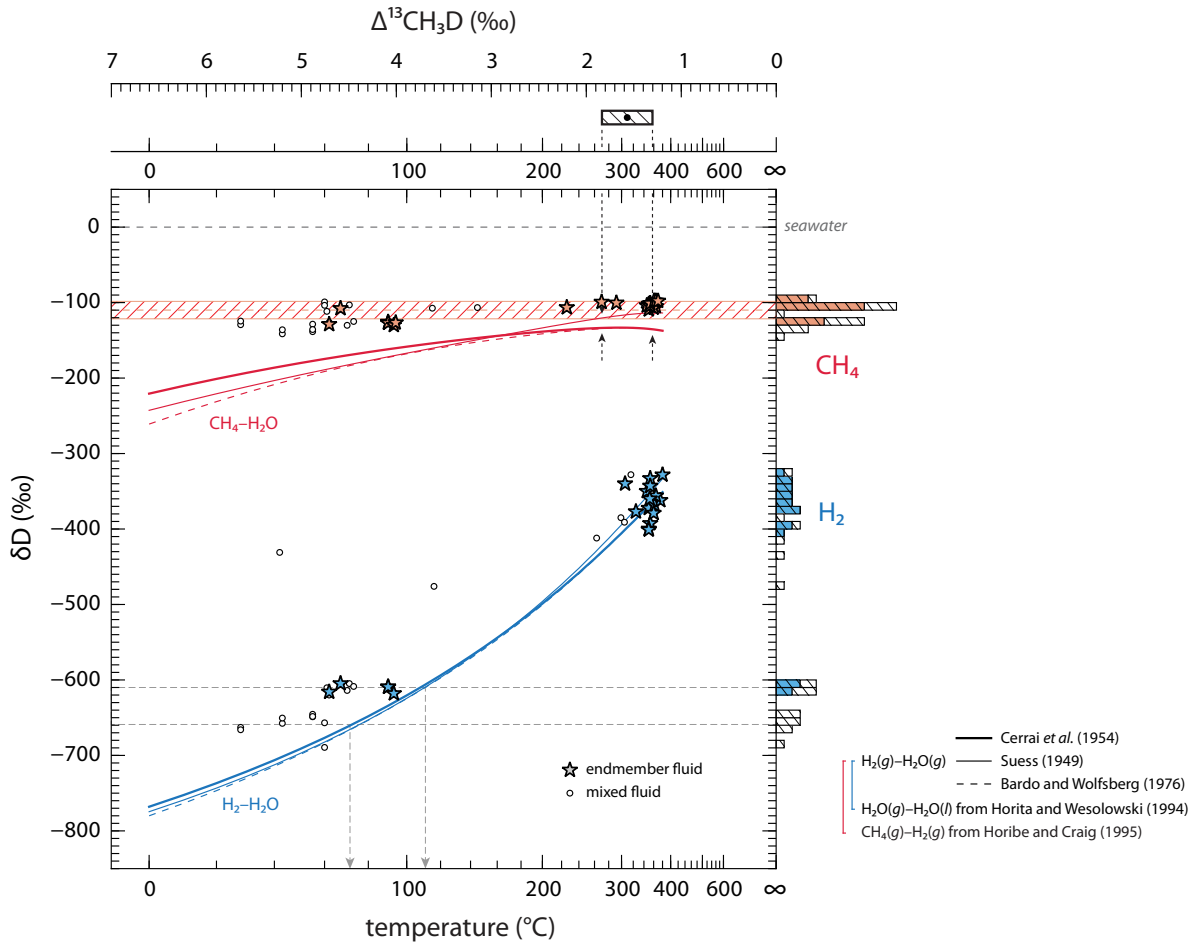
8. FIGURES





801

802 **Fig. 2.** Constraints on abiotic methane formation and stability from thermodynamics and clumped isotopologue
 803 data. **(A)** Plot of fugacity of H_2 as a function of temperature at 500 bar, after Shock (1992). The red line repre-
 804 sents the fugacity of H_2 at equilibrium, according to the reaction $CO_2(g) + 4H_2(g) \rightleftharpoons CH_4(g) + 2H_2O(l)$, when the
 805 fugacities of CH_4 and CO_2 are equal, and assuming unit activity for $H_2O(l)$. Grey lines represent equilibrium H_2
 806 fugacities buffered by the mineral assemblages hematite-magnetite (HM), pyrite-pyrrhotite-magnetite (PPM), fay-
 807 alite-magnetite-quartz (FMQ), and serpentine-magnetite-brucite (SBM). The curve for SBM is the is the low-Fe
 808 serpentinite from Sleep et al. (2004), and is truncated above 400 °C where serpentinization is unlikely to occur
 809 (see Sec. 4.2.3). Red shaded area represents the intersection of regions corresponding to geologically-relevant H_2
 810 fugacity and where CH_4 is thermodynamically stable relative to CO_2 . The black bar represents the temperature
 811 range over which our interpretation suggests that methane synthesis is both favorable and facile on timescales of
 812 relevance to hydrothermal systems. **(B)** Clumped isotopologue temperatures of methane from studied vents (data
 813 and error bars from Table 1). Equivalent $\Delta^{13}CH_3D$ values are plotted on the bottom axis, and are derived from
 814 the calibration of Wang et al. (2015). The dotted line and gray hatching represent the mean $\pm 1s$ of the $\Delta^{13}CH_3D$
 815 values across all studied vents ($1.57 \pm 0.28\%$, $n = 9$). The \times symbols mark measured vent temperatures (Table 2).



816

817 **Fig. 3.** Hydrogen isotopic composition of CH₄ (red) and H₂ (blue) in seafloor hydrothermal fluids plotted against
 818 measured vent temperatures. Data are from unconsolidated fields studied by Welhan and Craig (1983),
 819 Proskurowski et al. (2006), Kawagucci et al. (2010), Charlou et al. (2010), and us (see Tables 1 and 2), and are
 820 tabulated in Supplementary Table 2. Endmember fluids (identified by low Mg contents) are represented by stars,
 821 and vent fluids containing a mixture of hydrothermal endmember fluid and seawater are represented by circles.
 822 Data from sites exhibiting phase separation (Charlou et al., 2010) or with fluids diffusely effluxing through col-
 823 onies of deep-sea snails or shrimp (Kawagucci et al., 2010) are excluded from this plot (see note *f* under Supple-
 824 mentary Table 2). Red hatching indicates the average δD of CH₄ in endmember fluids ($-110 \pm 12\%$, 1s) in the
 825 compiled dataset. Red and blue curves represent the δD values of CH₄ and H₂ (respectively) in D/H equilibrium
 826 with seawater-like H₂O ($\delta D = 0\%$) calculated by combining published calibrations for H₂(g)/H₂O(g),
 827 H₂O(g)/H₂O(l), and CH₄(g)/H₂(g) (see Sec. 4.1). Note that measured values for δD of H₂O in fluids from Lost
 828 City are +2 to +7‰ (Proskurowski et al., 2006) and thus the equilibrium values for CH₄ and H₂ at Lost City are
 829 slightly (~5‰) higher than those indicated by the curves. Gray arrows near bottom of plot bracket the probable
 830 range of temperatures for closure of H₂-H₂O isotopic exchange. The bar at top represents the mean ± 1s of meas-
 831 ured $\Delta^{13}\text{CH}_3\text{D}$ values and corresponding clumped isotopologue temperatures (310^{+53}_{-42} °C) reported in Table 1,
 832 and black arrows point to the range of δD values of CH₄(g) in equilibrium with seawater at the temperatures indi-
 833 cated by $\Delta^{13}\text{CH}_3\text{D}$ data. At right is a histogram of the δD data.

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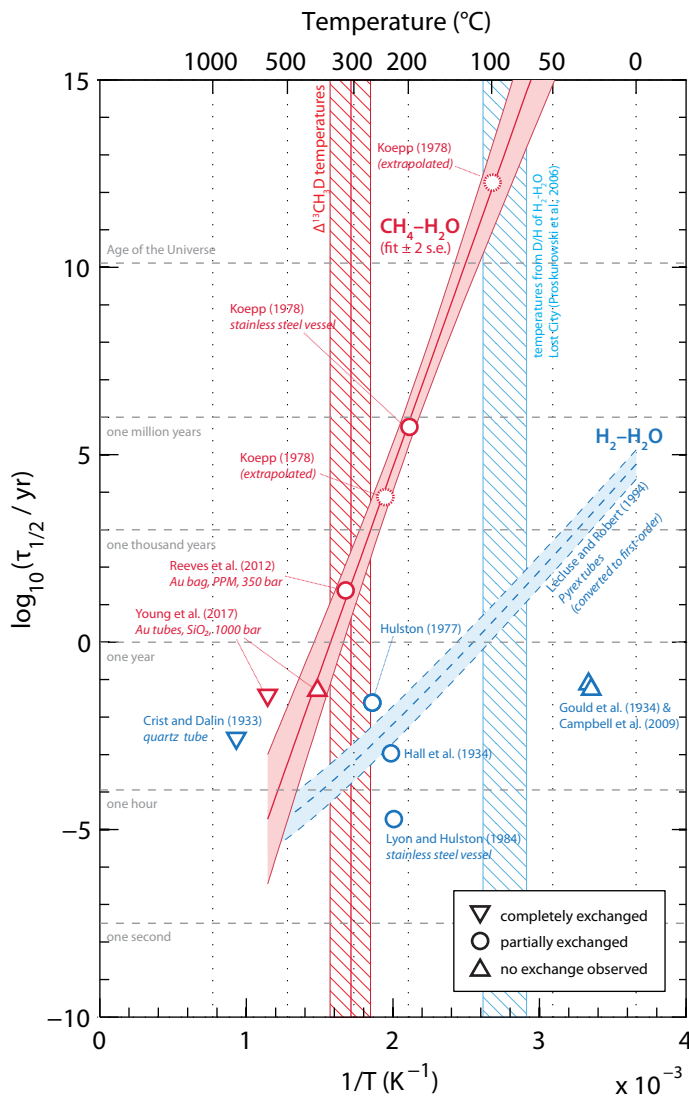
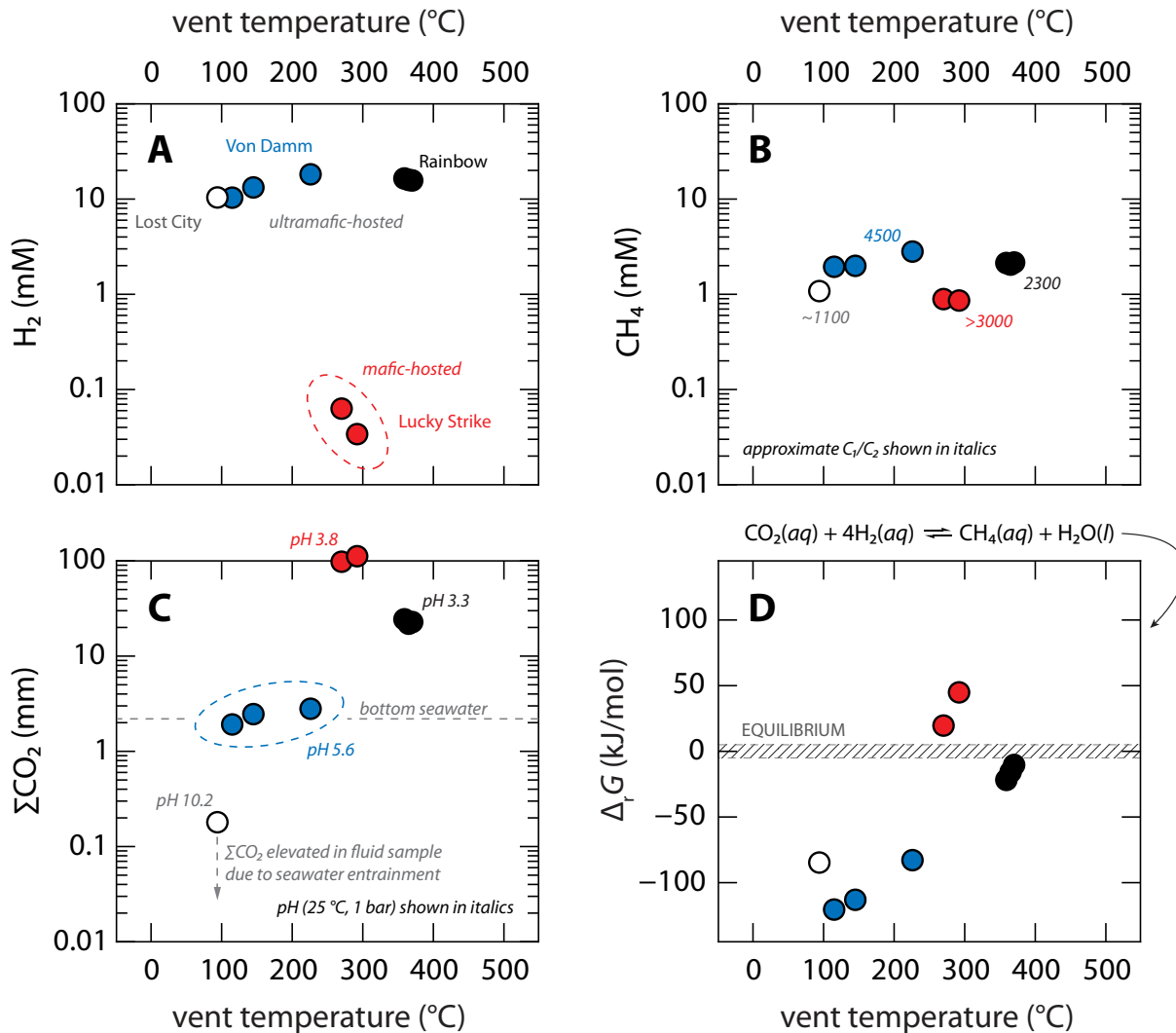
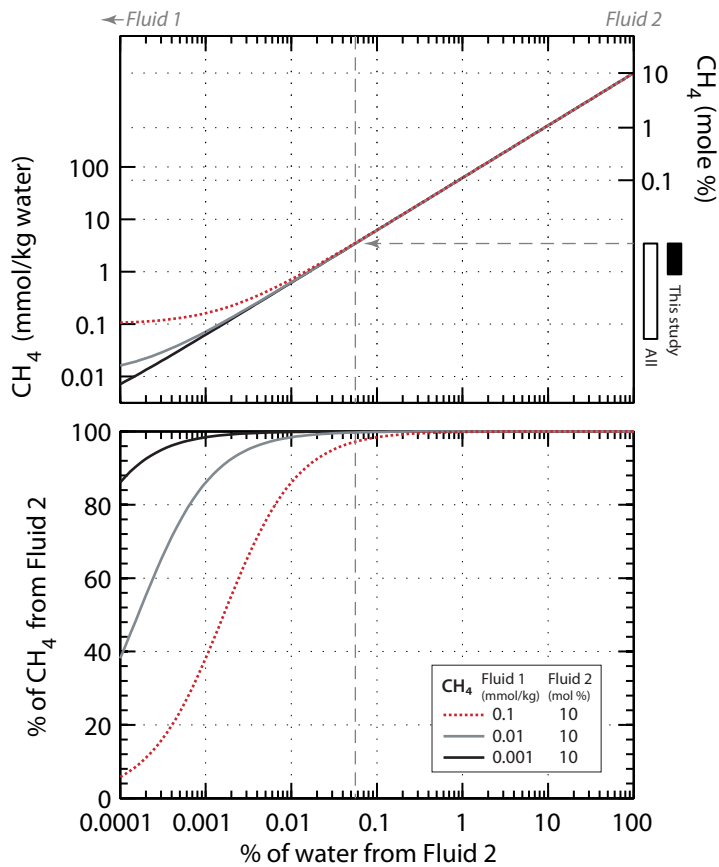


Fig. 4. Half-exchange timescales ($\tau_{1/2} = \ln(2) / k$) for hydrogen exchange between CH_4 & H_2O (red symbols) and H_2 & H_2O (blue) based on experiments done in the absence of added catalyst (Crist and Dalin, 1933; Hall et al., 1934; Gould et al., 1934; Hulston, 1977; Koepp, 1978; Lyon and Hulston, 1984; Lécluse and Robert, 1994; Campbell et al., 2009; Reeves et al., 2012; Young et al., 2017). Reactions were assumed to be first order in CH_4 or H_2 . For data points representing experiments for which rate constants were not reported or in which exchange was not observed, the y-axis value is the duration of the experiment. Downward- and upward-pointing triangles are, respectively, maximum and minimum estimates of the exchange timescale. The $\tau_{1/2}$ for CH_4 - H_2O exchange from Reeves et al. (2012) comes from Supplementary Fig. 1. Second-order rate coefficients for H_2 - H_2O exchange from Lécluse and Robert (1994) were converted to pseudo-first-order rate coefficients by multiplying by the vapor pressure of H_2O calculated at temperatures T and a pressure of 1 kbar (see Sec. 4.1). Uncertainties in exchange rates are difficult to estimate, but are probably several orders of magnitude. Clumped isotopologue temperatures for CH_4 from the present study (red hatched bar) and apparent temperatures from D/H geothermometry of H_2 - H_2O in endmember fluids at the Lost City site (blue hatched bar) (Proskurowski et al., 2006) are also shown.



879

880 **Fig. 6.** Composition of vent fluids and energetics of methane synthesis in aqueous phase. Concentrations of (A)
 881 H_2 , (B) CH_4 , and (C) ΣCO_2 are plotted against measured vent temperatures (data from Table 2). Also shown are
 882 molar ratios of methane to ethane (C_1/C_2) in (B), and pH values of endmember fluids in (C). (D) Gibbs energy of
 883 reaction for methane formation from CO_2 and H_2 in aqueous solution (Reaction 5), calculated at vent T and P
 884 conditions ($\Delta_r G$, Table 2). Gray hatching represents thermodynamic equilibrium (taken as $\Delta_r G = 0 \pm 5$ kJ/mol).
 885 Methane formation in aqueous solution is thermodynamically favorable for points plotting below the hatched area.
 886 Symbol colors are the same as those in Fig. 1.



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888 **Fig. 7.** Composition of fluids formed by mixing of a CH₄-poor actively-circulating seawater-derived hydrother-
 889 mal fluid (*Fluid 1*) with a CH₄-rich fluid such as those observed in inclusions in plutonic rocks on the Southwest
 890 Indian Ridge and on the Mid-Atlantic Ridge (*Fluid 2*) (Kelley, 1996; Kelley, 1997; Kelley and Früh-Green,
 891 1999). Mixing curves are plotted for CH₄ concentrations in the *Fluid 1* endmember ranging from 1 to
 892 100 μmol/kg (assumed values for background CH₄ that is e.g., possible to be derived from complete thermal alter-
 893 ation of dissolved organic matter in deep ocean water). Calculations assume that molalities of species other
 894 than CH₄ have a negligible effect on mole fractions in the high-CH₄ fluid. The black and white bars show CH₄
 895 concentrations in vent fluids from this study (Table 2) and from mid-ocean ridge hydrothermal systems globally
 896 (Keir, 2010).

897

9. TABLES

898

899 **Table 1**

900 Carbon and hydrogen isotope ratios and clumped isotopologue abundances of methane in studied hydrothermal fluids.

Field	Vent	Sample(s)	$\delta^{13}\text{C}$ (‰)	δD (‰)	$\Delta^{13}\text{CH}_3\text{D}$ (‰)	$T_{13\text{D}}$ (°C)
Rainbow	Guillaume	J2-352-IGT4	-17.6	-97.7	0.95 ± 0.60	450 +298/-136
	CMSP&P	J2-354-IGT3	-17.5	-97.8	1.50 ± 0.60	322 +142/-85
	Auberge	J2-352-IGT3	-17.4	-97.9	1.73 ± 0.60	285 +114/-73
Von Damm	Old Man Tree ^a	J2-612-IGT6/-IGT8	-16.2	-107.4	1.71 ± 0.35	288 +60/-47
	Ravelin 1	J2-617-IGT6	-16.4	-106.6	1.56 ± 0.60	312 +134/-82
	East Summit	J2-612-IGT2	-16.4	-106.5	1.35 ± 0.60	350 +167/-95
Lost City	Beehive	J2-361-IGT5/-CGTW _u	-10.9	-126.6	1.84 ± 0.60	270 +104/-68
Lucky Strike	Medea ^a	J2-359-IGT2/-CGTY	-14.2	-99.3	1.63 ± 0.40	301 +75/-55
	Isabel ^a	J2-357-IGT5/-CGTY	-12.6	-100.4	1.85 ± 0.30	269 +45/-37

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902 Values for $\delta^{13}\text{C}$, δD , and $\Delta^{13}\text{CH}_3\text{D}$ are reported relative to Vienna Pee Dee Belemnite (VPDB), Vienna Standard Mean Ocean
903 Water (VSMOW), and the stochastic distribution, respectively. Analytical uncertainties for $\delta^{13}\text{C}$ and δD are both ca. $\pm 0.1\text{‰}$
904 (95% confidence intervals). Uncertainties listed for $\Delta^{13}\text{CH}_3\text{D}$ and $T_{13\text{D}}$ are 95% confidence intervals; the last digit in each
905 (hundredths and ones places, respectively) is not significant.

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907 ^a Samples analyzed in duplicate. Uncertainties listed are 2 s.e.m. (standard error of the mean) of the replicate measurements
908 ($n = 2$).

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Table 2

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Fluid compositions^a used in thermodynamic calculations and calculated Gibbs energy of reaction ($\Delta_r G$) for abiotic methane formation via Reaction 5 at studied vent sites.^b

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Field	Vent	T (°C) ^c	P (bar)	pH ^d	ΣCO_2 (mm)	H_2 (mM)	CH_4 (mM)	$\Delta_r G$ (kJ/mol) ^e
Rainbow	Guillaume	361	230	3.33	24.3	16.5	2.13	-22
	CMSP&P	365	230	3.36	21.9	15.9	2.05	-16
	Auberge	370	230	3.35	22.8	15.7	2.16	-11
Von Damm	Old Man Tree ^f	115	235	5.81	1.80	10.5	1.97	-121
	Ravelin 1 ^f	145	235	5.83	2.52	13.4	2.02	-113
	East Summit	226	235	5.56	2.80	18.2	2.81	-83
Lost City	Beehive	96	70	10.20	0.18 ^g	10.4	1.08	-85
Lucky Strike	Medea	270	170	3.81	98.0	0.063	0.89	+20
	Isabel	292	170	3.81	112.0	0.034	0.86	+45

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Analytical uncertainties ($2s$) are ± 2 °C for T ; $\pm 5\%$ for H_2 , ΣCO_2 , and CH_4 ; and ± 0.05 units for pH. Abbreviations: mm, mmol/kg fluid; mM, mmol/L fluid.

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^a All concentrations shown are extrapolated to endmember fluid composition (regressed to zero Mg content), except where noted. Data are from McDermott et al. (2015) and Reeves et al. (2014).

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^b For each vent fluid, the energetic favorability of methane formation was assessed by calculating the Gibbs energy of reaction ($\Delta_r G$), defined by the relationship: $\Delta_r G = RT \ln(Q/K)$, where R is the universal gas constant, T is measured fluid temperature in kelvin, Q is the reaction quotient, and K is the equilibrium constant at T and seafloor pressure P . Equilibrium constants were calculated using thermodynamic data and standard states from Johnson et al. (1992) and Shock and Helgeson (1990). For all calculations, the activity of $\text{H}_2\text{O}(l)$ was assumed to be unity. Activity coefficients were assumed to be unity for neutral dissolved species. For all fluids except for that from Lost City,^g the concentration of $\text{CO}_2(aq)$ was assumed to be equal to ΣCO_2 , a reasonable approximation given the low measured shipboard pH values and calculated equilibrium speciation of dissolved carbonate species at in-situ temperatures and seafloor pressures.

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^c Maximum measured vent temperature.

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^d Shipboard pH measurement (25 °C and 1 atm).

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^e A negative value of $\Delta_r G$ indicates a thermodynamic drive for the reaction to proceed as written from left to right (i.e., methane formation favored). Given uncertainties associated with chemical analyses and thermodynamic data, calculated $\Delta_r G$ values within ± 5 kJ/mol of zero are interpreted to indicate that the reaction has approached or attained a state of thermodynamic equilibrium (Seewald, 2001).

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^f Concentrations for the fluids from Old Man Tree and Ravelin 1 vents at Von Damm were not extrapolated to zero Mg. Consistent high magnesium contents in duplicate gas-tight samples (14.0 and 15.0 mmol/kg fluid, respectively) indicate that fluids of the hot-spring source had mixed with seawater near the surface prior to discharge (McDermott et al., 2015), so they are not true zero-Mg endmembers. While CH_4 concentrations are lower in these fluids than in the Mg-deficient fluid from East Summit, ratios of CH_4 isotopologues did not change during subsurface mixing (Table 1).

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^g An arbitrary $\text{CO}_2(aq)$ concentration of 1 nmol/kg was used in thermodynamic calculations for the Lost City fluid, similar to Reeves et al. (2014). The actual concentration value is subject to substantial uncertainty due to difficulties in determining the near-zero endmember ΣCO_2 content in vent fluids, given that some entrainment of ΣCO_2 -replete seawater always occurs during sampling (Proskurowski et al., 2008). Varying this value by as much as ten orders of magnitude would not affect the conclusion that methane formation is thermodynamically favorable in the fluid, due to the high H_2 content and the power of 4 to which the activity of $\text{H}_2(aq)$ is raised in the mass action expression.

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Supplementary material for:

Clumped isotopologue constraints on the origin of methane at seafloor hot springs

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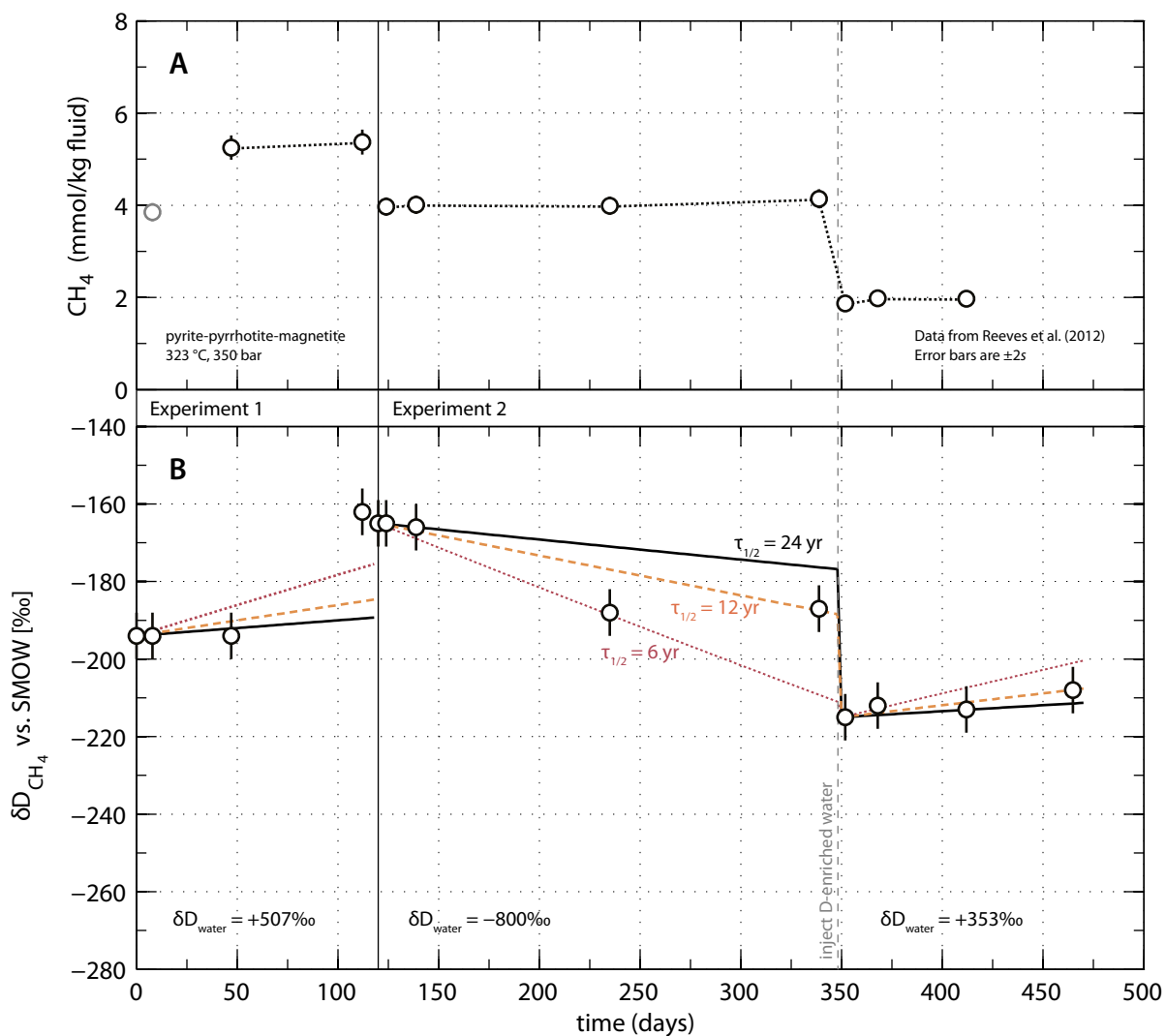
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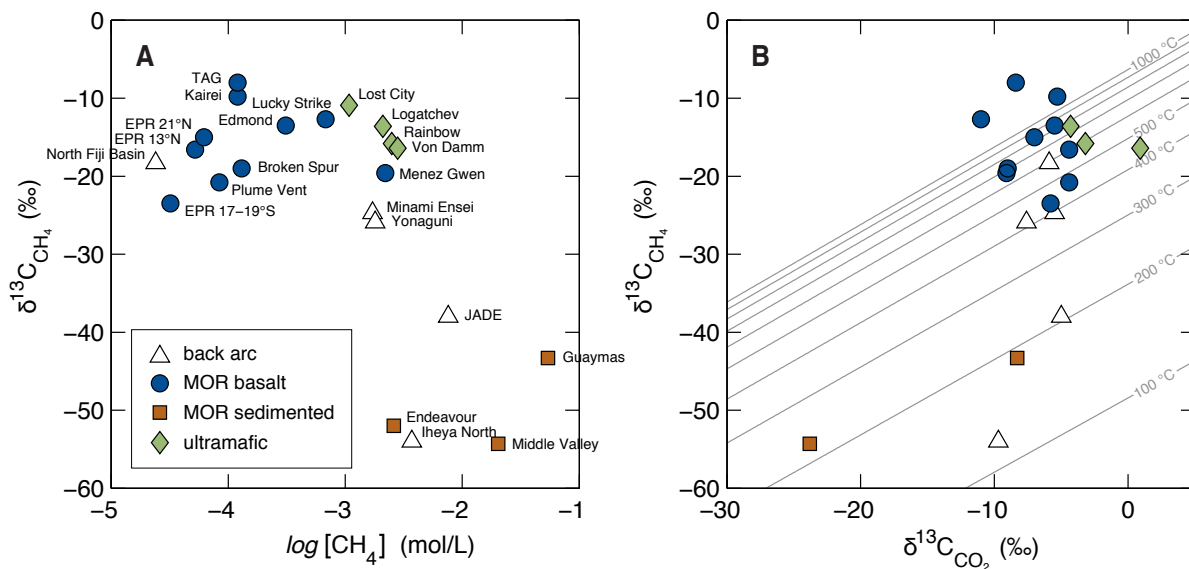
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Supplementary Fig. 1. Experimental constraints on hydrogen exchange between $\text{CH}_4(\text{aq})$ and $\text{H}_2\text{O}(\text{l})$ from two experiments conducted by Reeves et al. (2012) in a flexible cell hydrothermal apparatus (gold-titanium oxide) at 323 °C and 350 bar. Concentrations of CH_4 (A) remain indistinguishable within analytical error ($\pm 5\%$, $2s$) in Experiments 1 and 2. In Experiment 1, hydrocarbon concentrations in the first timepoint (8 days, shown in gray) were considered erroneously low, and can be excluded (see Reeves et al., 2012 for further details). Measured pH was ~ 4.2 , and concentrations of H_2 and $\Sigma\text{H}_2\text{S}$ were 0.26–0.7 mmol/kg fluid and ~ 11 mmol/kg fluid, respectively, consistent with predictions for a Fe–S–O–H aqueous fluid buffered by pyrite-pyrrhotite-magnetite (PPM) at experimental conditions (Reeves et al., 2012). Panel (B) shows measurements of D/H of CH_4 compared against modeled kinetics for D/H exchange with varying half-exchange time ($\tau_{1/2} = \ln(2) / k$). The modeled kinetics assume that CH_4 concentration is constant, the rate of isotopic exchange is first order in CH_4 , and the equilibrium D/H fractionation factor $[(\text{D}/\text{H})_{\text{methane}}/(\text{D}/\text{H})_{\text{water}} - 1]$ is -130‰ (see Figs. 3 and 5). We take $\tau_{1/2} = 24$ yr (black curve) as a best-guess estimate of the rate of true isotopic exchange; this value is shown in Fig. 4.



Supplementary Fig. 2. Carbon isotope data for CH_4 plotted against (A) CH_4 concentration and (B) $\delta^{13}\text{C}$ of inorganic carbon. Data from both sediment-influenced and unsedimented vent fields are shown. Data are from Charlou et al. (2000), Charlou et al. (2002), Cruse and Seewald (2006), Ishibashi et al. (1995), Charlou et al. (1996), Ishibashi et al. (1994), Kawagucci et al. (2011), Kumagai et al. (2008), Merlivat et al. (1987), Lein et al. (2000), Lilley et al. (1993), McDermott et al. (2015), Proskurowski et al. (2006), Proskurowski et al. (2008), Reeves et al. (2014), McCollom (2008), Konno et al. (2006), Welhan and Craig (1983), Welhan and Lupton (1987) and Evans et al. (1988). The curves in (B) are isotherms that represent carbon-isotopic equilibrium via the reaction: $^{12}\text{CH}_4(\text{g}) + ^{13}\text{CO}_2(\text{g}) \rightleftharpoons ^{13}\text{CH}_4(\text{g}) + ^{12}\text{CO}_2(\text{g})$ (Horita, 2001).

Supplementary Table 1

Measurements of methane samples tested for QA/QC. Each line in the table represents a $\sim 1 \text{ cm}^3$ SATP quantity of CH_4 that was purified and analyzed in the same manner as the vent fluid samples reported in Table 1.

Sample	Sample(s)	$\delta^{13}\text{C}$ (‰)	δD (‰)	$\Delta^{13}\text{CH}_3\text{D}$ (‰)	$T_{13\text{D}}$ (°C)
AL1 ^a	PRAC-1 + PRAC-5 ^c	-35.0	-149.9	2.38 \pm 0.29	207 +31/-27
	PRAC-2 + PRAC-3 ^c	-34.6	-147.1	2.16 \pm 0.40	230 +51/-40
	PRAC-4	-33.8	-146.6	2.26 \pm 0.26	220 +29/-25
NGS-3 ^b	Hydrothermal Test	-73.5	-176.2	5.33 \pm 0.49	39 +18/-16

Values for $\delta^{13}\text{C}$, δD , and $\Delta^{13}\text{CH}_3\text{D}$ are reported relative to Vienna Pee Dee Belemnite (VPDB), Vienna Standard Mean Ocean Water (VSMOW), and the stochastic distribution, respectively. Analytical uncertainties for $\delta^{13}\text{C}$ and δD are both ca. $\pm 0.1\text{‰}$ (95% confidence intervals). Uncertainties listed for $\Delta^{13}\text{CH}_3\text{D}$ and $T_{13\text{D}}$ are 95% confidence intervals; the last digit in each (hundredths and ones places, respectively) is not significant.

^a Accepted values of $\delta^{13}\text{C}$, δD , $\Delta^{13}\text{CH}_3\text{D}$, and $T_{13\text{D}}$ of AL1 are -34.5‰ , -147.7‰ , $+2.41 \pm 0.07\text{‰}$, and $204 \pm 7 \text{ °C}$ (95% confidence intervals), respectively (Wang, 2017).

^b NGS-3 was determined in triplicate by Wang et al. (2015). The average of their reported values for $\delta^{13}\text{C}$, δD , $\Delta^{13}\text{CH}_3\text{D}$, and $T_{13\text{D}}$ are $-72.84 \pm 0.08\text{‰}$, $-176.04 \pm 0.23\text{‰}$, $+5.13 \pm 0.06\text{‰}$, and $46 \pm 2 \text{ °C}$ (95% confidence intervals), respectively.

^c These are small aliquots of AL1 (~ 0.2 to 0.8 cm^3 SATP CH_4) that were prepared in separate serum bottles and pooled during purification prior to analysis. This was done to verify that our pooling procedure maintains isotopic integrity.

Supplementary Table 2

Compilation of hydrogen isotope ratios of CH₄ and H₂ and associated data on vent fluids from unsedimented hydrothermal systems.

Field	Vent ^a	T_{\max} (°C) ^b	Mg (mM) ^c	ΣCO_2 (mm)	H ₂ (mM)	CH ₄ (mM)	$\delta^{13}\text{C}$ (‰)		δD (‰)			Notes	
							ΣCO_2	CH ₄	CH ₄	H ₂	H ₂ O ^d		
<i>Mid-Atlantic Ridge</i>													
Rainbow	Guillaume (X4)	361	0*	24.3	16.5	2.13	—	-17.6	-98	—	—	(1, 2)	
	CMSP&P	365	0*	21.9	15.9	2.05	—	-17.5	-98	—	—	(1, 2)	
	Auberge (X3)	370	0*	22.8	15.7	2.16	—	-17.4	-98	—	—	(1, 2)	
	—	365	0*	16	16	2.5	-3.2 ^e	-17.7	-105	-356	—	(3)	
	—	360	0*	17	13	1.6	-2.5 ^e	-17.8	-107	-379	—	(3)	
Lost City	Beehive	94	0*	0.18	10.4	1.08	—	-10.9	-127	—	—	(1, 2)	
		90	0*	—	—	—	—	—	-127	-609	+2 to 7	(4)	
	90	0*	—	—	—	—	—	—	-126	-609	+2 to 7	(4)	
	Marker 6	67	0*	—	—	—	—	—	—	-108	-605	+2 to 7	(4); cf. T_{\max} 96 °C in ref. 2
		62	0*	—	—	—	—	—	—	-129	-616	+2 to 7	(4); cf. T_{\max} 96 °C in ref. 2
	IMAX (IF)	55	—	—	—	—	—	—	—	-129	-649	+2 to 7	(4)
		55	—	—	—	—	—	—	—	-139	-646	+2 to 7	(4)
		55	—	—	—	—	—	—	—	-136	-648	+2 to 7	(4)
	Marker 7	28	—	—	—	—	—	—	—	-129	-663	+2 to 7	(4)
		28	—	—	—	—	—	—	—	-125	-666	+2 to 7	(4)
	Marker 8	43	—	—	—	—	—	—	—	-141	-658	+2 to 7	(4)
		43	—	—	—	—	—	—	—	-136	-651	+2 to 7	(4)
	Marker C	62	—	—	—	—	—	—	—	-126	-620	+2 to 7	(4)
		70	—	—	—	—	—	—	—	-130	-614	+2 to 7	(4)
	Marker H	60	—	—	—	—	—	—	—	-99	-657	+2 to 7	(4)
		60	—	—	—	—	—	—	—	-104	-689	+2 to 7	(4)
	Marker 3	61	—	—	—	—	—	—	—	-112	-610	+2 to 7	(4)
		71	—	—	—	—	—	—	—	-103	-605	+2 to 7	(4)
		73	—	—	—	—	—	—	—	-125	-609	+2 to 7	(4)
	—	93	0*	—	—	—	—	-11.9	-130	-618	—	(3)	
Broken Spur	—	353	0*	—	—	—	—	—	—	-393	—	(4)	
Logatchev	—	350	0*	—	—	—	—	—	-109	-372	—	(4); Logatchev 1?	
Logatchev 1	—	346	0*	3.6	9	2.0	+4.1 ^e	-10.2	-104	-350	—	(3)	
	—	352	0*	4.4	13	2.6	+7.4 ^e	-10.3	-104	-360	—	(3)	
Logatchev 2	—	320	0*	6.2	11	1.2	+9.5 ^e	-6.1	-93	-231	—	(3); phase-separated ^f	
Ashadze 1	—	353	0*	3.7	8	0.5	+2.1 ^e	-12.3	-104	-333	—	(3)	
	—	353	0*	—	19	1.2	+4.6 ^e	-14.1	-101	-343	—	(3)	
Ashadze 2	—	296	0*	—	26	0.8	+0.2 ^e	-8.7	-107	-270	—	(3); phase-separated ^f	
Lucky Strike	Medea	270	0*	98	0.063	0.89	—	-14.2	-99	—	—	(1, 2)	
	Isabel	292	0*	112	0.034	0.86	—	-12.6	-100	—	—	(1, 2)	

Field	Vent ^a	T_{\max} (°C) ^b	Mg (mM) ^c	ΣCO_2 (mm)	H_2 (mM)	CH_4 (mM)	$\delta^{13}\text{C}$ (‰)		δD (‰)			Notes
							ΣCO_2	CH_4	CH_4	H_2	$\text{H}_2\text{O}^{\text{d}}$	
<i>East Pacific Rise</i>												
9° N	—	380	0*	—	—	—	—	—	—	-328	—	(4)
21° N	Nat. Geo. Soc.	350	0*	—	30.5	1.4	-7.0	-15.0	-102	-401	+0.5	(5, 6)
<i>Central Indian Ridge</i>												
Kairei	Kali	362	0*	8.0	3.3	0.12	-5.3	-9.8	—	-368	—	(7, 8)
		316	8.4	12.1	3.6	—	—	—	—	-328	—	(7, 8)
	Monju	299	5.2	7.9	2.1	—	—	—	—	-385	—	(7, 8)
		42	50.9	9.3	8×10^{-4}	—	—	—	—	-431	—	(7, 8)
		87	43.7	6.0	0.69	—	—	—	—	-361	—	(7, 8); snail colony ^f
	Fugen	22	48.3	12.6	0.13	—	—	—	—	-493	—	(7, 8); snail colony ^f
		305	4.5	9.5	2.7	—	—	—	—	-391	—	(7, 8)
	Daikoku (Marker 30)	306	0*	—	2.2	—	—	—	—	-340	—	(7)
		—	350	0*	—	—	—	—	—	-400	—	(4)
	Edmond	Nura Nura	375	0*	12.8	0.11	0.31	-5.5	-13.5	—	-362	—
Marker 27		325	0*	12.3	0.10	—	—	—	—	-377	—	(7, 8)
White Head		263	12.4	8.1	0.04	—	—	—	—	-412	—	(7, 8)
Grand Shrimp Valley		281	13.4	12.1	0.48	—	—	—	—	-681	—	(7, 8); shrimp colony ^f
Marker 24		116	40.6	8.7	0.07	—	—	—	—	-476	—	(7, 8)
<i>Mid-Cayman Rise</i>												
Von Damm	Old Man Tree	115	14.0	1.80	10.5	1.97	—	-16.2	-107	—	—	(1, 9)
	Ravelin 1	145	15.0	2.52	13.4	2.02	—	-16.4	-107	—	—	(1, 9)
	East Summit	226	0*	2.80	18.2	2.81	—	-16.4	-107	—	—	(1, 9)

Data sources: (1) this study; (2) Reeves et al. (2014); (3) Charlou et al. (2010); (4) Proskurowski et al. (2006); (5) Welhan and Craig (1983); (6) Horibe and Craig (1995); (7) Kawagucci et al. (2010); (8) Kumagai et al. (2008); (9) McDermott et al. (2015).

Abbreviations: mm, mmol/kg fluid; mM, mmol/L fluid.

^a Dash (—) indicates that data were not reported or that samples were unable to be matched across multiple references.

^b Maximum measured vent temperature.

^c Asterisk (*) indicates near-endmember fluid sample (represented by stars in Fig. 4). For these samples, concentrations of ΣCO_2 , H_2 , and CH_4 and $\delta^{13}\text{C}$ values of ΣCO_2 have been extrapolated to endmember fluid composition (regressed to zero Mg content) assuming entrainment of seawater containing ~53 mM Mg.

^d Endmember vent fluids typically have δD values of H_2O between -2 and +4‰ (Shanks et al., 1995). A value of 0‰ was assumed when no data could be found (see text and Fig. 4).

^e Values are as reported; it is not known whether correction for ΣCO_2 in seawater was applied.

^f Figure 4 excludes data from these fluids, which have either effluxed through macrofaunal colonies or are venting with atypically low salinity.

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