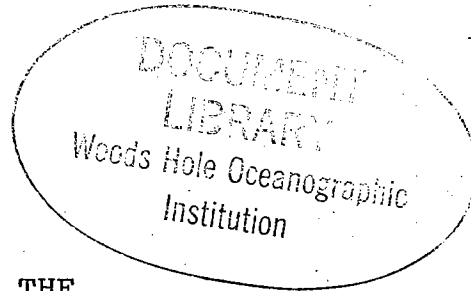


THE DISTRIBUTION OF DISSOLVED SILICA IN THE  
DEEP WESTERN NORTH ATLANTIC OCEAN

by

Gerald J. Needell

B.S., Northeastern University  
(1975)



SUBMITTED IN PARTIAL FULFILLMENT OF THE  
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November, 1978

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To Sally - who is my inspiration

## ABSTRACT

The distribution of dissolved silica in the deep western North Atlantic Ocean is presented. The potential temperature-dissolved silica relationship is compared with the potential temperature salinity relationship in the North Atlantic Deep Water. Geographical variations in the potential temperature-dissolved silica relationship are discussed with particular emphasis on the low silica signal of the Western Boundary Undercurrent (WBUC). The WBUC is shown to have a significant influence on the potential temperature-dissolved silica relationship from the tail of the Grand Banks of Newfoundland to Cape Hatteras.

It is suggested that a region of enhanced mixing is present west of 65°W that is responsible for the observed changes in the dissolved silica distribution.

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Submitted to the Department of Meteorology  
on November 13, 1978 in partial fulfillment of the requirements  
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Thesis Supervisor: Dr. J. R. Luyten, Associate Scientist

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## I. Introduction

The distribution of dissolved silica in the deep western North Atlantic Ocean may be used to identify water masses with distinct characteristics within the potential temperature and salinity range of the ubiquitous North Atlantic deep water.

The most prominent feature of the silica distribution is a narrow band of low silica along the continental rise. This low silica feature shall be identified as the Western Boundary Undercurrent (WBUC) in keeping with the literature (Richardson (1977)). It shall be demonstrated that the WBUC can be traced from the tail of the Grand Banks of Newfoundland to Cape Hatteras. With the data presented below, a "global" view of the WBUC shall be provided.

In the past, observations of the WBUC have consisted of hydrographic sections across the continental rise for which the geostrophic transport was computed. Either neutrally buoyant floats or moored current meters were used to set absolute velocities (Swallow and Worthington (1961), Volkmann (1962), Richardson (1977), Clarke, Hill, Reiniger, and Warren (1978)). Occasionally, dissolved silica was measured and it could be demonstrated that the water observed did have a northern origin (Richardson (1977), Clarke, Hill, Reiniger, and Warren (1978)).

Eight month long current meter records along 70°W show a steady westward flow on the upper continental rise (Luyten,

1977). Unfortunately, one cannot be sure whether the westward flow is the WBUC (i.e., low silica) or is part of a broad westward flow in the slope water region (Webster (1969)).

This brings up an important problem with the historical observations of WBUC. Stommel (1957) postulated the existence of a deep southward flowing western boundary current in the Atlantic Ocean. His model was based on the hypothesis that some thermohaline process caused sinking of water in arctic regions and upwelling in antarctic regions. To conserve mass, he argued that a deep flow would be present along the western boundary.

In the following year, Swallow and Worthington (1961) occupied hydrographic sections and tracked neutrally buoyant floats near Cape Romaine, North Carolina. In the hydrographic data, they found that the deep isotherms sloped downward off shore indicating the presence of westward shear. When they observed that the neutrally buoyant floats moved westward, it was assumed that the WBUC had been found. The current meter data from 70°W shows that even where the flow is relatively steady toward the west, it is not always so. Occasional reversals in the flow are not uncommon. Thus one must be cautious in the interpretation of few day-long float tracks.

It should also be noted that the spatial coverage of the data from which we can observe the WBUC is not uniform. The only long-term current meters on the continental rise were at 70°W. Hydrographic sections with some method of determining absolute velocities are only available near 70°W and Cape



Hatteras. The results of a section at  $50^{\circ}\text{W}$  are not considered realistic for reasons discussed in the text. Thus we are left with only a limited set of observations of the WBUC. There is no direct evidence for a continuously flowing current that passes around the tail of the Grand Banks of Newfoundland and follows the continental rise toward the west. To the contrary, the current meter data and the variations in the transport calculations suggest that whatever flow there is, is not steady.

The distribution of dissolved silica will be used to demonstrate that the WBUC can be identified in terms of its water mass characteristics. The effects of the WBUC can be traced along the continental rise from the tail of the Grand Banks of Newfoundland to Cape Hatteras. However, the structure of the flow cannot be determined from the available data.

Evidence is presented below which suggests that a region of enhanced mixing is present west of  $65^{\circ}\text{W}$  and has a pronounced effect upon the dissolved silica distribution.

## II. Background

### 2.1 Historical uses of dissolved silica as a water mass tracer

It has only been in the past three decades that one has been able to map out the distribution of dissolved silica. On the basis of sparse data, Sverdrup, Johnson, and Fleming (1942) concluded that the concentration of dissolved silica

increased monotonically with depth everywhere in the world oceans. They pointed out that the silica values in the Atlantic Ocean were in general, lower than those of the other oceans. Cooper (1952) examined the sources and sinks of dissolved silica and suggested that the concentration of silica in the deep water should be in equilibrium and might be useful as a water mass tracer. Without any observations to rely on, Cooper assumed that the tundra drainage in the Arctic would result in a high silica content of the waters of the North Polar Sea. He then reasoned that there would be a flow of high silica water in the Greenland and Irminger Currents which could be easily detected by its silica content. While recent observations have indeed shown that there exists high silica water in the Baffin Basin ( $\text{SiO}_2 > 80 \mu\text{g A/l}$  - Corwin and McGill (1963), Grant (1970)), there is no evidence of it getting over the sill and into the Labrador Basin. In fact, the Norwegian and Labrador Seas are characterized by a low silica signal ( $\text{SiO}_2 \sim 10-15 \mu\text{g A/l}$ ). As it shall be demonstrated, these waters are easily detected along the western boundary of the western North Atlantic Ocean.

Armstrong (1965) showed the distinctly lower concentrations of silica in the deep Atlantic Ocean, when compared to the world oceans. He concluded that this was a result of the unique exchange of bottom waters that occurs in the Atlantic Ocean and suggested that "...our understanding of water movements and of mixing processes in the Atlantic Ocean would be greatly increased by more knowledge of the silica content of Arctic waters."

Metcalf (1969) showed that the concentration of dissolved silica could be used to identify Antarctic Bottom Water, North Atlantic Deep Water, and Antarctic Intermediate Water. He also suggested that silica might be useful in studying the NADW where little variation in the potential temperature-salinity relationship is found.

Carmack (1973) found that silica was a valuable tracer in the Antarctic. He found that he could distinguish between water masses that appeared indistinguishable on the basis of  $\theta$ -S data alone.

## 2.2 The deep western North Atlantic Ocean

It is well known that there is a tight correlation between potential temperature and salinity in the deep western North Atlantic Ocean for potential temperatures below  $3.0^{\circ}\text{C}$  (Worthington and Metcalf (1961), Wright and Worthington (1970)). Most of the water characterized by this  $\theta$ -S correlation for  $\theta > 1.9^{\circ}\text{C}$  have been classified as North Atlantic Deep Water (NADW) (Wright and Worthington (1970)). There is little geographic variation in the  $\theta$ -S distribution of the western North Atlantic in this potential temperature range (Worthington and Wright (1970)). This has made it difficult for physical oceanographers to study the deep general circulation in terms of potential temperature and salinity characteristics.

Near the continental rise, geostrophic calculations and direct velocity measurements (neutrally buoyant float and

moored current meters) allowed observers to deduce the existence of a deep southward flowing current (Swallow and Worthington (1961); Volkmann (1962); Richardson (1977)). This flow has become known as the Western Boundary Undercurrent and is responsible for the equatorward movement of deep water from the Norwegian and Labrador Seas (Stommel (1957); Richardson (1977); Clarke, Reiniger, Hill, and Warren (1978)).

Metcalf (1969) demonstrated the usefulness of the dissolved silica content of deep water as a water mass tracer. Using the available data, he constructed a map of the concentration of dissolved silica on the 2.0°C potential temperature surface. His figure is reproduced as Fig. 1. The principal feature of this map is the general north-south gradient of dissolved silica. The paucity of data prevents a detailed analysis of the distribution but this map does suggest that dissolved silica is a useful tracer of the general circulation of the North Atlantic Ocean.

Richardson (1977) and Worthington (personal communication) found that there is a drop in the silica concentration along potential isotherms as one approaches the continental rise.

Worthington (1976) has proposed a model of the general circulation of the North Atlantic Ocean, in which he postulated the existence of two anti-cyclonic gyres. One to the north of the Grand Banks of Newfoundland, and the second comprised of the Gulf Stream and its return flow bounded to

the south by the Mediterranean salt tongue. In the deep water Worthington based his interpretation partly on the change in the dissolved silica distribution in the deep water across the proposed boundary between the two gyres. His figure showing characteristic deep  $\theta$ -SiO<sub>2</sub> plots for the two gyres is reproduced as Fig. 2. Worthington's model also allows for a WBUC passing around the tail of the Grand Banks of Newfoundland and along the continental rise until it crosses beneath the Gulf Stream near Cape Hatteras. He allows for no exchange of water between the WBUC and the southern gyre.

Clarke, Hill, Reiniger, and Warren (1978) have criticized Worthington's model. They claim that in the deep water, no evidence exists for two distinct gyres, and that the silica distribution (from a survey of their own) indicates that there is an exchange of water between the "Northern Gyre" area and the Sargasso Sea.

In this work the distribution of dissolved silica in the deep western North Atlantic Ocean, based on several sections occupied during the last two decades, shall be presented. A large portion of the data is from the last six years. After presenting a summary of the data and methods of analysis used, we shall describe the water mass characteristics present and form a consistent picture of the geographical variations in the silica distribution. It shall be demonstrated that there are significant variations of the silica distribution within the NADW. These variations

will allow us to trace the origins of some components of the NADW which cannot be distinguished on the basis of  $\theta$ -S characteristics. The distribution of silica on potential temperature surfaces shall be examined and these distributions compared with models of the general circulation. In addition the WBUC shall be examined by following its silica signal along the continental rise from the tail of the Grand Banks to Cape Hatteras. Evidence shall be presented for isopycnal mixing along the WBUC path, resulting in downstream dilution of the low silica WBUC water.

The data coverage of the western North Atlantic Ocean has expanded greatly over that which was available to Metcalf (1969) and Worthington and Wright (1972). More structure in the silica distribution can be discerned here than was possible previously. The time span over which the data was collected must certainly be kept in mind. Little is known of the temporal variations of the deep water properties. That a consistent picture of the deep silica distribution can be presented is taken to indicate that these temporal variations are indeed small.

It must, of course, be kept in mind that local anomalies do exist. They may be either due to temporal, spatial, or possibly instrumental effects. As discussed in Appendix I large anomalies which contradict surrounding data must be suspect, but more reasonable local features are retained in the data set and are accepted as essential characteristics of the ocean. Where these features stand out they will be noted.

### III. Data and Methods

The data used in this report was obtained by searching the WHOI and NODC oceanographic files for cruises in which temperature, salinity, and dissolved silica were measured. Only those cruises where salinities were measured by a conductivity bridge (i.e., three decimal place precision) were used. This basically restricted the data set to cruises since 1958. The data used is summarized in Table 1 (Fig. 3). The data from one cruise, KNORR 12, was rejected for reasons given in Appendix I.

In a few cases, duplicate sections (i.e., two or more sections at the same location) existed in the data. The section with more closely spaced stations was used. In addition, cruises for which samples were analyzed at sea were preferred to those during which the samples were frozen. All of the data used in this report was analyzed at sea with either a spectrophotometer or autoanalyzer.

For all cruises except AII 100 (see below) the data were accepted as they existed in the files. Observations denoted as questionable by the observers were discarded. Plots of  $\theta$ -SiO<sub>2</sub> and S-SiO<sub>2</sub> were formed and observations which disagreed with the main body of data were rejected (see discussion of  $\theta$ -SiO<sub>2</sub> data). All of the data (again with the exception of AII 100 as noted below) were from previously published sources (see Table I) and had been edited on the basis of the mean  $\theta$ -S distribution of the North Atlantic Ocean.

Table I.

Year	Ship	Cruise	Location	Previous Publication
1958	DISCOVERY II	I.G.Y. 2	24° 30' N* 32° 15' N*	Metcalf (1969) Worthington and Wright (1970)
1961	ATLANTIS	265	40° 15' N*	
1961	CHAIN	20	Woods Hole-Bermuda	Worthington and Wright (1970)
1966	ATLANTIS II	18	near 36°N,66°W and 36°N,63°w	Metcalf(1969)
1971	TRIDENT		Cape Hatteras	Richardson (1977)
1972	HUDSON		50°W and around the tail of the Grand Banks	Clarke, Hill, Reiniger and Warren (1978)
1972	CHAIN	104	around the tail of the Grand Banks	Clarke, Hill, Reiniger and Warren (1978)
1972	KNORR	30	GEOSECS stations* 27-34	Broecker, Takahashi and Li(1976)
1975	KNORR	48	64° 30'W	Worthington(in preparation)
1977	KNORR	66	55°W	Worthington(in preparation)
1978	ATLANTIS II	100	70°W	

\* used only in the preparation of the maps of dissolved silica on potential temperature surfaces



The section from cruise AII 100 was occupied by the author in the spring of 1978. The samples were collected with a 24-bottle rosette of Niskin bottles on a Neil Brown CTD. For this work only the temperature data from the CTD was used. The salinities were determined at sea on a Guildline Autosal salinometer and the dissolved silica was determined with a spectrophotometer. The data were edited by plotting both salinity and silica versus pressure and removing those points from which both salinity and silica values indicated contamination of the sample. For the thirteen stations occupied, the worst had eight "contaminated" samples and the best had none (there were 24 samples per station). The final  $\theta$ -S data set was then compared to the Worthington-Metcalf (1969) mean  $\theta$ -S curve and no other questionable salinities were obvious below 4.0°C.

The contour plots presented in Section V were prepared by linearly interpolating the  $\theta$ -SiO<sub>2</sub> data for each station to the  $\theta$  values used.

#### IV. Water Mass Characteristics

##### 4.1 Potential temperature salinity characteristics

As previously stated, the deep western North Atlantic Ocean is characterized by a tight potential temperature salinity correlation. Figure 4 is a composite plot of deep  $\theta$ -S observations for those stations indicated. The boundaries of the North Atlantic Deep Water are indicated (Worthington

and Wright (1970)). Below the inflection point at  $\theta \sim 1.85^\circ\text{C}$  a low salinity tail is present indicating the effects of Antarctic Bottom Water (Wright and Worthington (1970); Broecker, Takahashi, and Li (1976)).

#### 4.2 Potential temperature-silica characteristics

Figure 5 is a composite plot of deep observations of potential temperature and dissolved silica for those stations indicated. It should be noted that within the potential temperature range of the North Atlantic Deep Water ( $\theta > 1.9^\circ\text{C}$ ), we find the maximum scatter in the  $\theta$ - $\text{SiO}_2$  plot.

Note also that for  $\theta < 1.85^\circ\text{C}$  there is a sharply increasing silica tail corresponding to the low salinity tail of the Antarctic Bottom Water (AABW) (Broecker, Takahashi, and Li (1976)). This high silica branch is confined to the south of the tail of the Grand Banks of Newfoundland. One can see the difference clearly when the data are separated into groups of data from either side of the tail of the Grand Banks (see Figures 6a and 6b). Those stations east of  $50^\circ\text{W}$  show no AABW. This geographical limit to the northward influence of the AABW is similar to that used by Wright and Worthington (1970) in their volumetric potential temperature, salinity census of the North Atlantic Ocean. The AABW was seen clearly in the North American Basin, but not in the Labrador Basin (Figures 7a and 7b were reproduced from Wright and Worthington (1970)).

Wright and Worthington (1970) also found that in the Labrador Basin (but absent in the North American Basin) there

was a relatively saline (with respect to the Antarctic Bottom Water) and cold ( $\theta < 1.9^{\circ}\text{C}$ ) water mass (Denmark Straits, Overflow Water (DSOW)). In agreement with this we find that for  $\theta < 1.9^{\circ}\text{C}$ , east of  $50^{\circ}\text{W}$  (Figure 6b), there is a low silica branch on the  $\theta$ - $\text{SiO}_2$  plot. This water is not generally observed west of  $50^{\circ}\text{W}$  (see exception noted below).

The two branches of the  $\theta$ - $\text{SiO}_2$  plot come together at  $\theta \sim 1.9^{\circ}\text{C}$  as did the two branches of the  $\theta$ -S distribution (Wright and Worthington (1970)). Worthington and Wright (1970) showed that at  $1.8^{\circ}\text{C}$  and below there is no direct path between the AABW and the DSOW. It is between  $1.8^{\circ}\text{C}$  and  $1.9^{\circ}\text{C}$  a direct path first exists (Figures 8a and 8b).

The data presented above suggest that for  $\theta < 1.9^{\circ}\text{C}$  both salinity and dissolved silica can be used to identify the AABW and the DSOW.

In the range of the NADW ( $\theta > 1.9^{\circ}\text{C}$ ) one may take advantage of the wide variation of dissolved silica (whereas there is very little variation of salinity) to distinguish between water masses of different origin within the NADW.

In the data from east of  $50^{\circ}\text{W}$  (Figure 6b) there is a distinct gap in the data from  $\theta = 1.9^{\circ}$  to  $3.4^{\circ}\text{C}$ . The higher silica branch runs from  $\theta \sim 3.4^{\circ}\text{C}$  and  $\text{SiO}_2 \sim 18 \mu\text{g A/l}$  to  $\theta \sim 1.9^{\circ}\text{C}$  and  $\text{SiO}_2 \sim 32 \mu\text{g A/l}$  where it joins in with the AABW and the DSOW. In this same potential temperature range, there exists a lower silica branch running from  $\theta \sim 3.4^{\circ}\text{C}$ ,  $\text{SiO}_2 \sim 16 \mu\text{g A/l}$  to  $\theta \sim 1.9^{\circ}\text{C}$ ,  $\text{SiO}_2 \sim 18 \mu\text{g A/l}$ . This lower silica water may be traced back into the Labrador and

Norwegian Seas (Grant (1970)), and is presumed to be the source of the Western Boundary Undercurrent (WBUC) (Richardson (1976); Clarke, Hill, Reiniger, and Warren (1978)).

West of 50°W, the distribution is quite different (Figure 6a) in the NADW range of potential temperature. There are only a few observations to the low side of the majority of the observations. The distinct gap seen to the east is not evident here. In fact, most of the "anomalous" observations come from two cruises. The section along 50°W shows the two branched distributions clearly and also shows some of the DSOW water at  $\theta \sim 1.8^\circ\text{C}$  (Figure 9). This section was grouped with the western group because of the AABW seen for  $\theta < 1.9^\circ\text{C}$  and  $\text{SiO}_2 > 30 \mu\text{g A/l}$ . The other source of the anomalous data is AII 18 (Figure 10). These data also show evidence of DSOW ( $\theta \sim 1.85^\circ\text{C}$ ,  $\text{SiO}_2 \sim 24 \mu\text{g A/l}$ ). These are the only observations of the water mass at a point so far west. Since these observations of the DSOW are so isolated, it is assumed that there was an intrusion of this water mass into the area at the time these data were collected.

With the two sections discussed above deleted, the  $\theta - \text{SiO}_2$  distribution west of 50°W consists of a broad line of values from  $\theta = 3.5^\circ\text{C}$  and  $\text{SiO}_2 = 14 \mu\text{g A/l}$  to  $\theta = 1.9^\circ\text{C}$   $\text{SiO}_2 = 30 \mu\text{g A/l}$  where it meets the AABW (Figure 11).

We stated previously that it was, in general, the lower values of silica that were found closer inshore. This will become more clear when the charts of silica on constant

potential temperature surfaces are discussed in Section V. The reason for the lower values inshore is assumed to be the influence of the WBUC (Richardson (197~~3~~); Worthington (1976)).

Examination of the individual sections which extend from the abyssal plain up onto the continental rise reveals that while there is a general tendency toward lower values of silica inshore, the strength of the gradient and the extreme values vary considerably (Figures 12a-e). The figures may be divided into two categories. The first are those sections which show a distinct gap between the WBUC and the remainder of the NADW (HUDSON, KN 48). The second are those which show a continuous band of observations but are limited to a range of values within the extremes of HUDSON and KN 48 (AII 100, CH 20, TRIDENT). These types of distributions are grouped geographically, the two branched distributions being to the east (east of 65°W) and the more uniform distributions being to the west.

#### 4.2.1 Polynomial fits to the potential temperature silica relationship

To quantify the differences between these distributions, polynomial fits to the data were made. It was found that little was gained for fits of higher than fourth order. The fits were obtained with a least square polynomial regression program for the Hewlett-Packard 9830 calculator. The fits and residuals are presented in Figures 13-17. Root mean square deviations for each of the fits were also computed

(Table 2). It should be noted that the r.m.s. deviations are lower to the west than to the east, the plots of the residuals show that for CH 20, AII 100, and TRIDENT, the residuals are evenly distributed between  $\pm 2\sigma$  with an increase of the scatter toward lower potential temperatures. For both KN 48 and HUDSON, we find that the residuals are evenly distributed above  $\theta = 3.0^\circ\text{C}$  but break into two branches for  $\theta < 3.0^\circ\text{C}$ . The low residuals can readily be identified in the WBUC water which forms the low silica branch of the  $\theta - \text{SiO}_2$  plots. To illustrate the effects of the WBUC on the fits, the inshore stations were removed (containing the WBUC observations) from the KN 48 and HUDSON sections and re-computed the fits (Figures 18, 19). The r.m.s. deviations are now 1.10 for HUDSON and 1.01 for KN 48. The plots of the residuals now show a more uniform distribution than did the fits including the WBUC. The curves fitted to the KN 48 ( $65^\circ\text{W}$ ) and HUDSON ( $50^\circ\text{W}$ ) sections with the WBUC removed are nearly identical (Figure 20). To the west of  $65^\circ\text{W}$ , the three fits produce similar curves. Above  $2.0^\circ\text{C}$  those for TRIDENT and AII 100 are nearly identical and CHAIN 20 shows a deviation from the other two between  $2.5^\circ$  and  $4^\circ\text{C}$  (Figure 21). In Figure 21 a 2 standard deviation envelope has been included to show that the discrepancies are not statistically significant even from  $2.5^\circ\text{C} \rightarrow 4^\circ\text{C}$ .

To demonstrate that the WBUC has an effect on the distribution of silica west of  $65^\circ\text{W}$  the fits for AII 100, CH 20,

Table II.

Cruise	Location	Standard Deviation (ugA/l)
HUDSON	50°W	2.72
KNORR 48	65°W	1.46
CHAIN 20	Woods Hole-Bermuda	1.41
ATLANTISII 100	70°W	1.10
TRIDENT	Cape Hatteras	1.21

and TRIDENT were plotted along with HUDSON and KN 48 (With WBUC removed) (Figure 22). The  $2\sigma$  envelope is from KN 48. It can be shown that for  $2.00 < \theta < 2.40^\circ\text{C}$  there is a significant difference between the AII, CH 20, TRIDENT curves and the HUDSON and KN 48, indicating that the water of the WBUC has a significant effect west of  $65^\circ\text{W}$ .

It should be emphasized that the data being compared were collected over a span of several years. The three sections which show the effect of the WBUC west of  $65^\circ\text{W}$  were occupied in 1961, 1971, and 1978 (CH 20, TRIDENT and AII 100, respectively). The consistency of the data suggests that the conditions were not anomalous at the time of the observations.

For comparison, the  $\theta \sim \text{SiO}_2$  plots for the sections discussed above are presented again with the KN 48 (without WBUC) curve and  $2\sigma$  envelope (Figures 23a-h).

Thus it has been shown while there are dramatic changes in the  $\theta - \text{SiO}_2$  distribution from the tail of the Grand Banks of Newfoundland to Cape Hatteras, one can demonstrate that the low silica signal of the WBUC can be traced throughout the region.

#### V. The Distribution of Dissolved Silica on Surface of Constant Potential Temperature

In the previous section, the plots of silica versus potential temperature demonstrated the existence of geographic variations in the silica distribution of the deep western North Atlantic Ocean. In particular, it was found that near



the tail of the Grand Banks of Newfoundland and along the continental rise as far west as 65°W, the North Atlantic deep water consisted of two distinct water masses, easily distinguished by their silica content. Further west, the water masses could not be so easily separated on the  $\theta$ -SiO<sub>2</sub> plots, however, it was shown that the low silica water did have a significant effect on  $\theta$ -SiO<sub>2</sub> relationship as far west as Cape Hatteras.

To illustrate further the geographic variability of the silica distribution, four potential temperature\* surfaces were chosen upon which to prepare contour maps of the silica distribution (Figures 24-27).

The data base used in preparing these maps is presented in Table I.

For potential temperatures less than 1.9°C one cannot trace water continuously throughout the North Atlantic Ocean. Figures 8a,b (Worthington and Wright (1970)).

Between 1.8°C and 1.9°C the bottom waters from northern and southern origin come in direct contact.

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\*Due to the tight  $\theta$ -S relationship in the deep western North Atlantic Ocean for  $\theta < 3.0^\circ\text{C}$  (Wright and Worthington (1970); Worthington and Metcalf (1961)) potential temperature surfaces nearly coincide with potential density surfaces.

### 5.1 The 1.9° potential temperature surface

The distributions of dissolved silica on the 1.9°C surface (Figure 24) shows the general southward increase in silica noted by Metcalf (1969) on the 2.0°C surface (see Figure 1).

Along the mid-Atlantic Ridge, a tongue of high silica water is found to extend northward. This is AABW and was identified by Metcalf (1969) using the same IGY sections.

In the vicinity of the tail of the Grand Banks of Newfoundland, there is a sharp silica gradient, with the lowest values inshore, following the bottom topography. One should note that this sharp gradient cannot be traced westward around the tail of the Grand Banks. The lowest contours ( $<22.5 \mu\text{g A/l}$ ) are only resolved by a few observations and our interpretation of this as a continuous feature may not be justified. Since none of the very low silica water penetrates around the tail of the Grand Banks, it is not associated with the WBUC, but rather with the Denmark Straits overflow water that is usually observed to be confined to the Labrador Basin (Figure 2, Wright and Worthington (1970); Worthington (1976)). Away from the continental rise the silica values are uniformly high ( $>35 \mu\text{g A/l}$ ) and join with the plume of AABW extending north. It should be noted that 1.9°C is the potential temperature at which the  $\theta\text{-SiO}_2$  and  $\theta\text{-S}$  (Wright and Worthington (1970)) plots show the joining of the DSOW and AABW (Figure 5).

West of 50°W, the silica distribution takes on a different character. There is a decrease in the silica values as

one approaches the continental rise but the sharp gradient of the DSOW is not present. After passing around the tail of the Grand Banks, the 30 and 32.5  $\mu\text{g A/l}$  contours diverge leaving a large area with little change in the silica concentration. This is the location of the deep anti-cyclonic gyre proposed by Worthington (1976).

It should also be noted that the 30  $\mu\text{g A/l}$  contour follows the topography of the continental rise until it disappears west of 70°W. It is commonly observed that the lowest silica contours will follow the topography until they end by running into the coast indicating some mixing process (Figures 24-27).

## 5.2 The 2.2°C potential temperature surface

On the 2.2°C potential temperature surface (Figure 25) the general north-south gradient is again present with a zonal gradient with the higher values of silica being found along the mid-Atlantic Ridge. There is a sharp gradient near the tail of the Grand Banks with the lowest value nearest the continental rise. This gradient may be traced along the continental rise and westward around the tail of the Grand Banks.

In contrast to the 1.9°C surface, the sharp gradient along the continental rise may be traced further west to 65°W. As it was previously defined, this is the WBUC transporting water southward from the deep Labrador Sea (Grant (1970)).

West of 65°W the gradient weakens. The 20  $\mu\text{g A/l}$  contour runs into the coast and the 22.5 and 25  $\mu\text{g A/l}$  contours diverge. Similar to the 1.9°C surface, the silica contours on the 2.2°C surface, away from the continental rise, near the Grand Banks diverge as one proceeds westward. Here the 25 and 27.5  $\mu\text{g A/l}$  contours diverge. The waters above the abyssal plain are characterized by a nearly uniform distribution of silica.

### 5.3 The 2.4°C potential temperature surface

The 2.4°C potential temperature surface, Figure 26, is similar to the 2.2°C surface. We see once again the northwest-southeast gradient with the highest values of silica found along the mid-Atlantic Ridge and the lowest along the continental rise.

One can trace a sharp silica gradient from the tail of the Grand Banks westward to 65°W where the 20 and 22.5  $\mu\text{g A/l}$  contours then diverge in a similar fashion to the 22.5 and 25  $\mu\text{g A/l}$  contours on the 2.2°C surface. Also similar to the 2.2°C surface, the 22.5 and 25  $\mu\text{g A/l}$  are close together near the tail of the Grand Banks but diverge when traced westward beyond 50°W. Thus it is found that the waters above the abyssal plain nearly uniform in silica.

The inshore gradient on the 2.4°C surface is not as sharp as that seen on the 2.2°C surface. As shown on the  $\theta\text{-SiO}_2$  plots (Figure 23a for example) the spread between the WBUC water and the remainder of the NADW increases with decreasing potential temperature.

#### 5.4 The 3.0°C potential temperature surface

The 3.0° potential temperature surface (Figure 27) shows a departure from the contribution on the previous surfaces. There is the expected southward increase in silica but at this potential temperature, there is open communication across the mid-Atlantic Ridge. The data indicate a sharp gradient near 32°N with abrupt changes above the mid-Atlantic Ridge. These features are not well resolved and shall only be pointed out here without further comment.

Near the tail of the Grand Banks, and along the continental rise, the sharp gradients seen on the previous surfaces are no longer present. The distribution is nearly uniform throughout the entire region north of 35°N from the Grand Banks to Cape Hatteras. This was shown also in the  $\theta$ -SiO<sub>2</sub> plots (see Figure 5) by the narrowing of the  $\theta$ -SiO<sub>2</sub> scatter for  $\theta \geq 3.0^\circ\text{C}$ .

## VI. Discussion

### 6.1 The $\theta$ -SiO<sub>2</sub> relationship

In the preceding sections a silica minimum (on surfaces of constant potential temperature) was found along the continental rise. This feature has been associated with the WBUC (Richardson (1977); Clark, Hill, Reiniger, and Warren (1978)). It has been demonstrated that the low silica water can be traced from the tail of the Grand Banks of Newfoundland to Cape Hatteras. However, it was also found